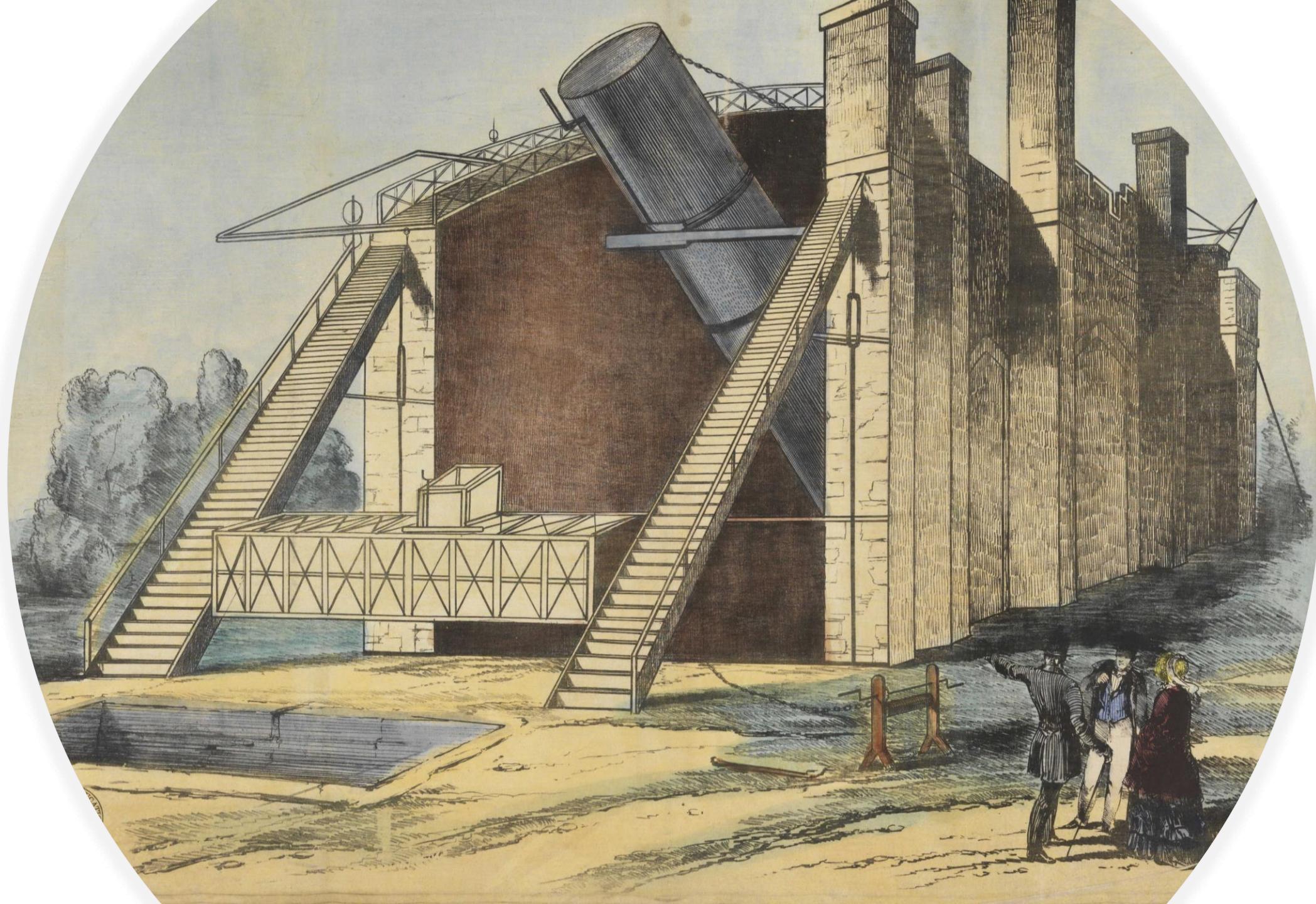




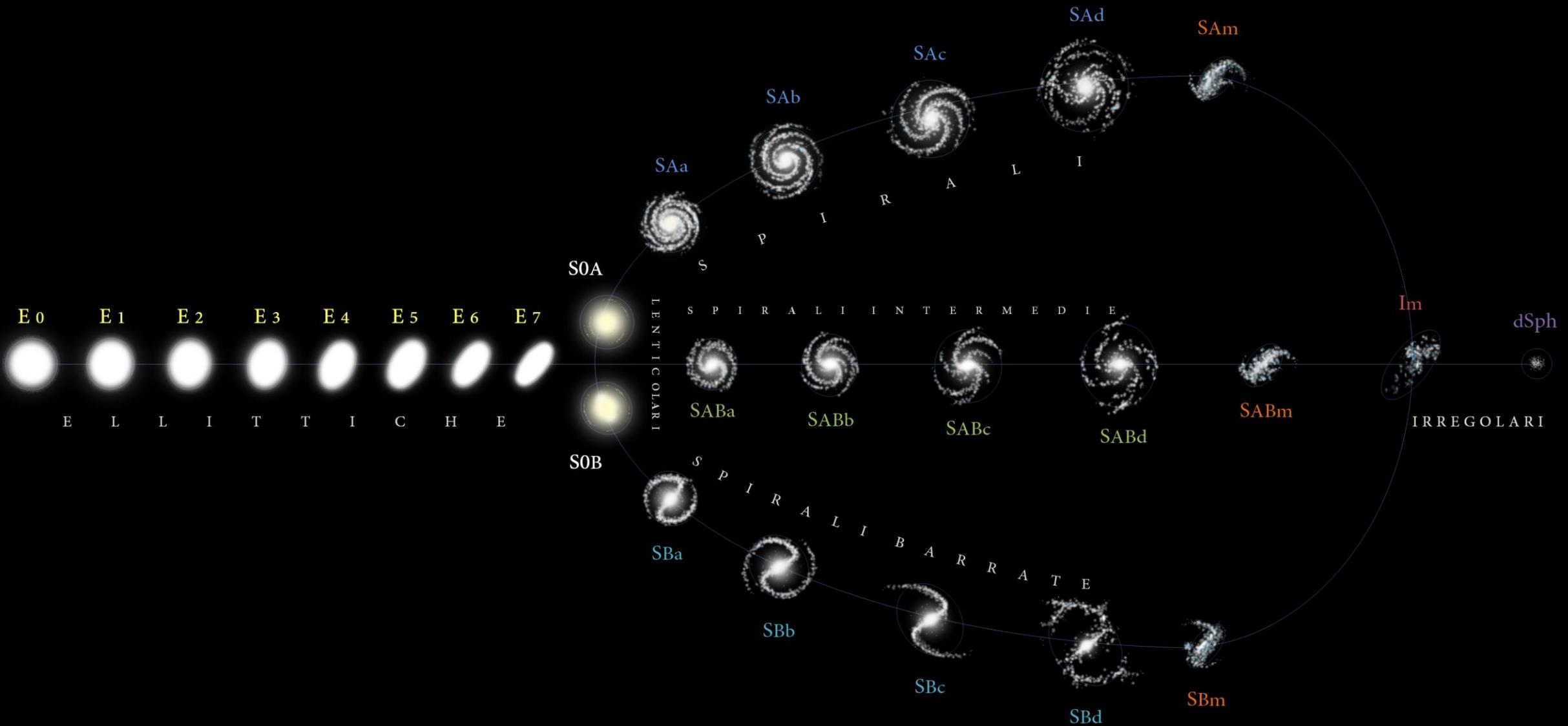
Dinamica delle galassie

L'importanza dell'ISM e delle stelle





Schema Hubble – Vaucouleurs



Radio

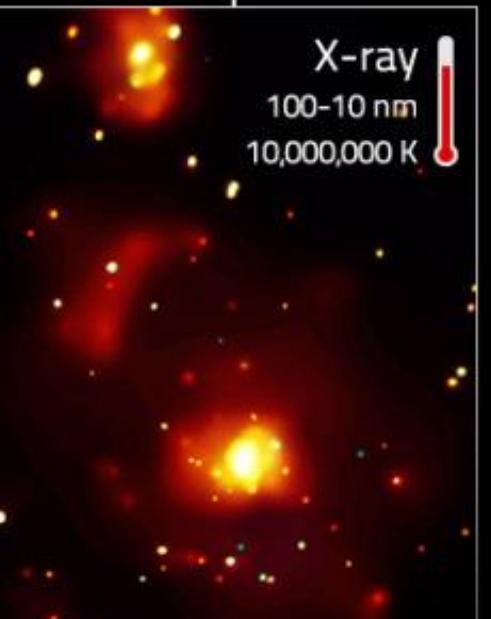
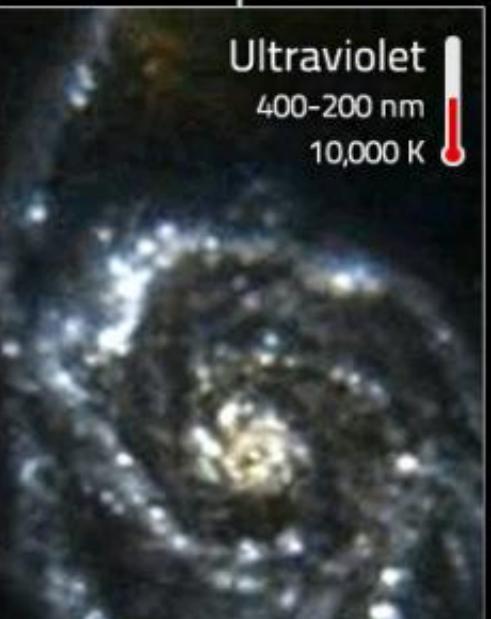
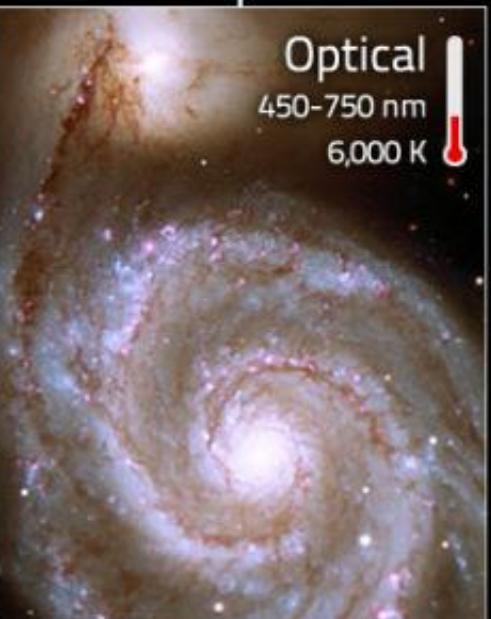
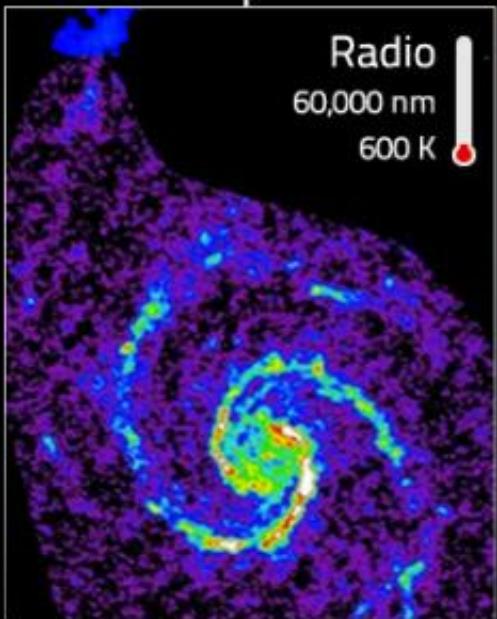
Microwave

Infrared

UV

X-Ray

Gamma Ray



Multiwavelength Whirlpool Galaxy

COLD GAS: Radio waves reveal regions of gas cool enough for CO₂ molecules to exist.

COOL STARS: Infrared shows smaller cool red stars that make up most of the galaxy.

SOLAR STARS: Optical light comes from stars around the size of the Sun.

HOT STARS: Ultraviolet shows the larger hot blue stars that are less frequent in galaxies.

HOT GAS: X-rays are emitted from the hottest regions of gas where atoms are ionized.

← COOL LOW ENERGY RADIATION

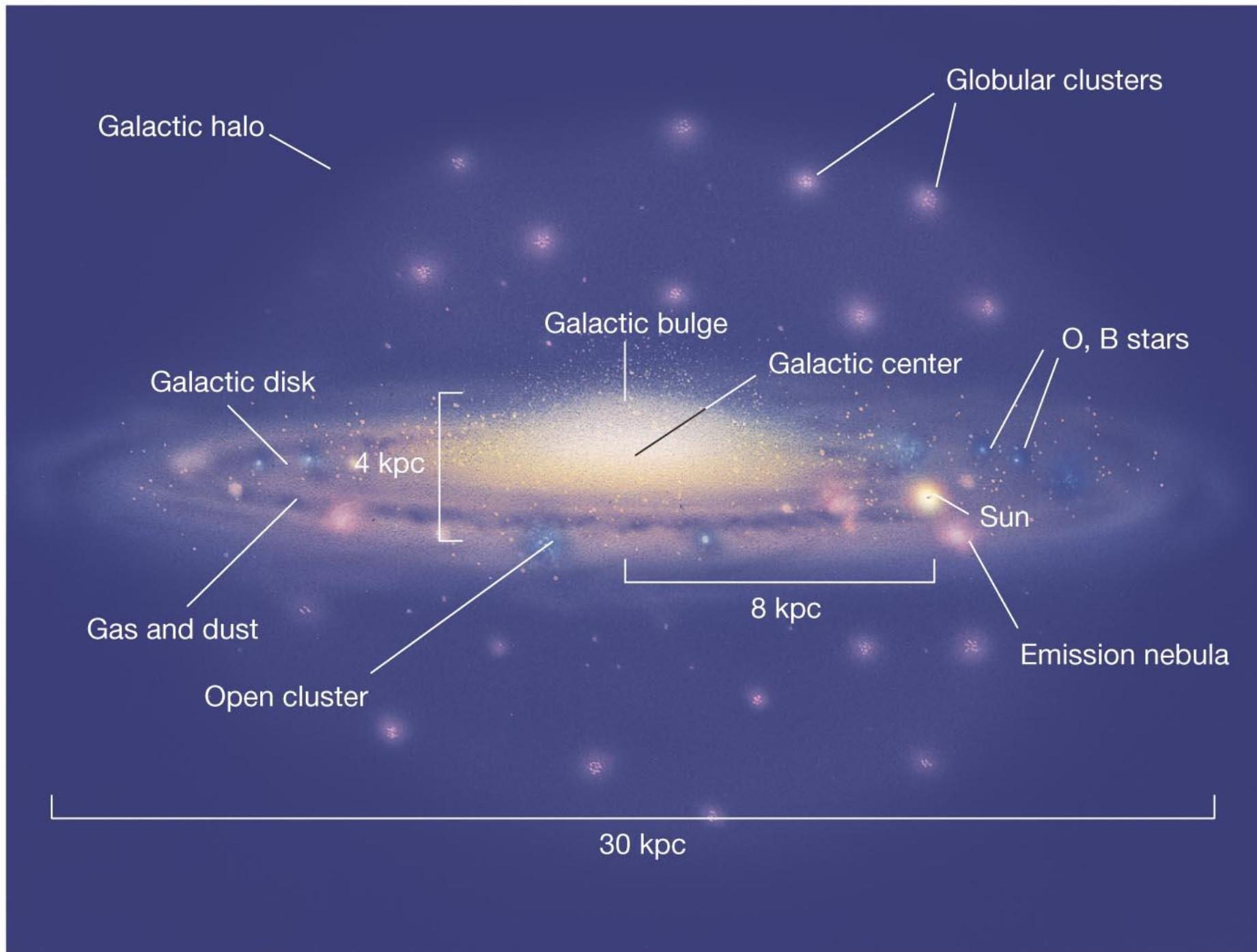
VISIBLE LIGHT

HOT HIGH ENERGY RADIATION →



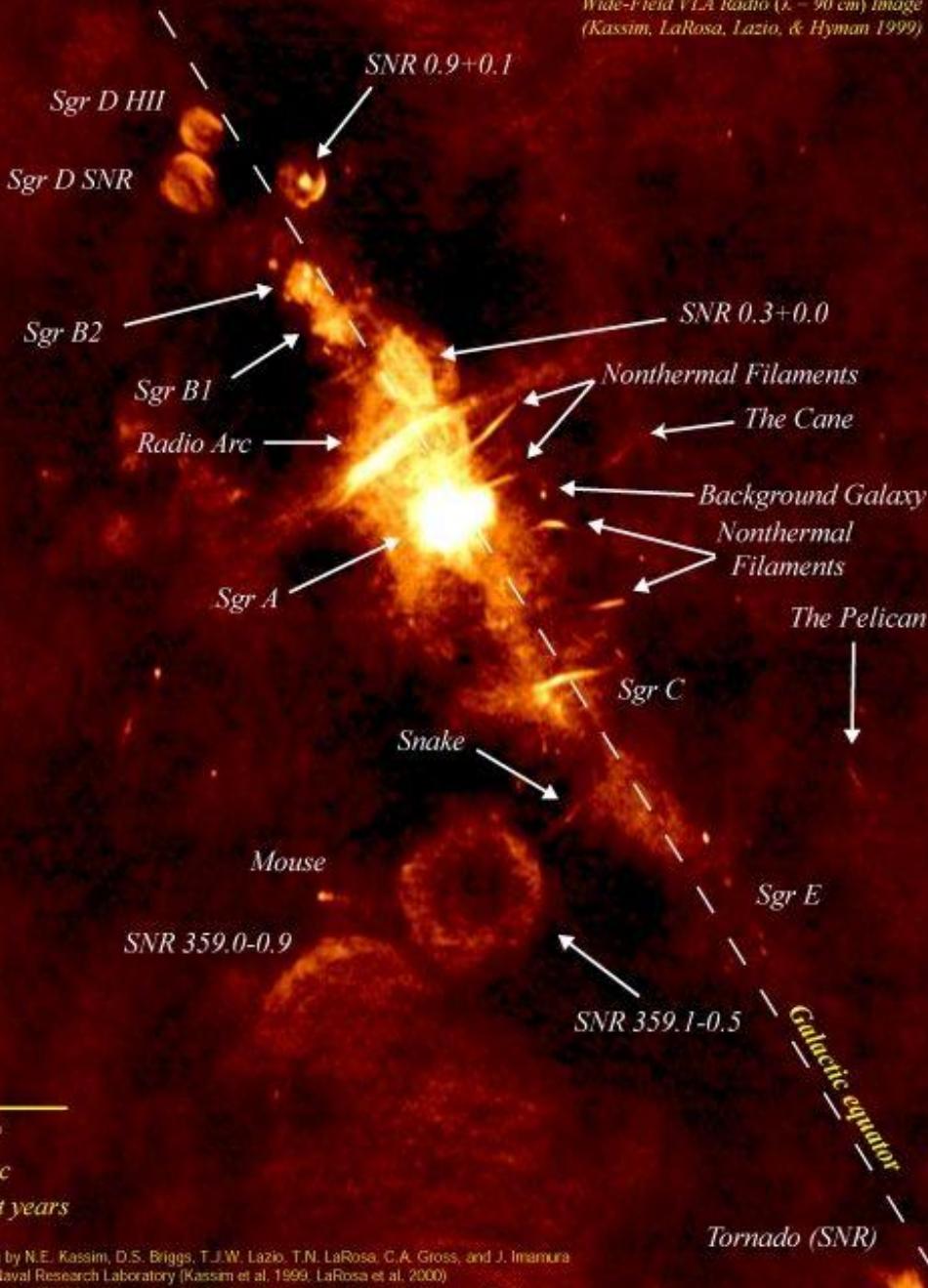
WAS A FIDGET SPINNER

BEFORE IT WAS COOL



The Galactic Center

Wide-Field VLA Radio ($\lambda = 90$ cm) Image
(Kassim, LaRosa, Lazio, & Hyman 1999)



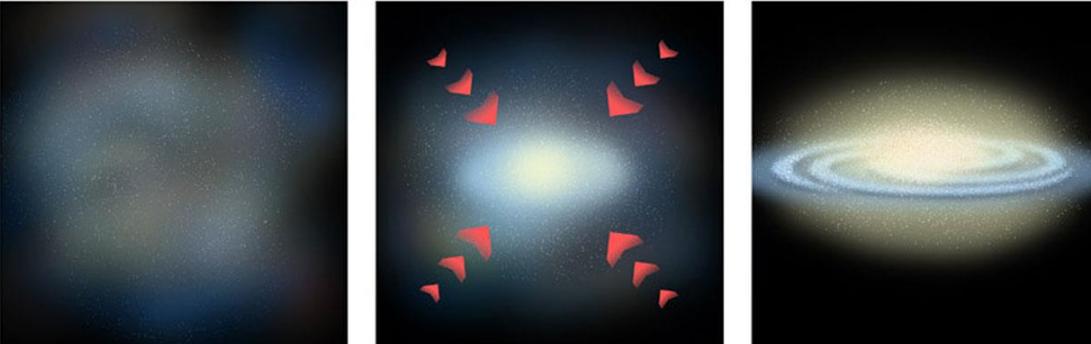
1992



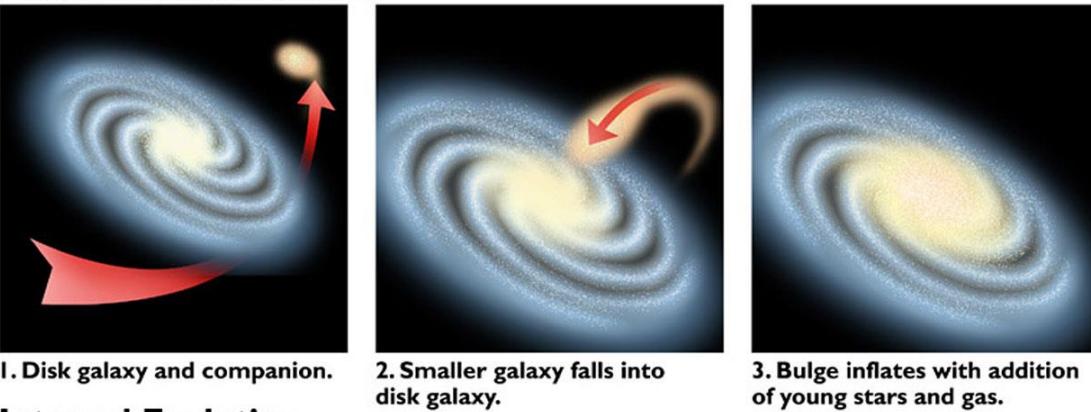
The Growth of Bulges in Spiral Galaxies

Three evolutionary scenarios

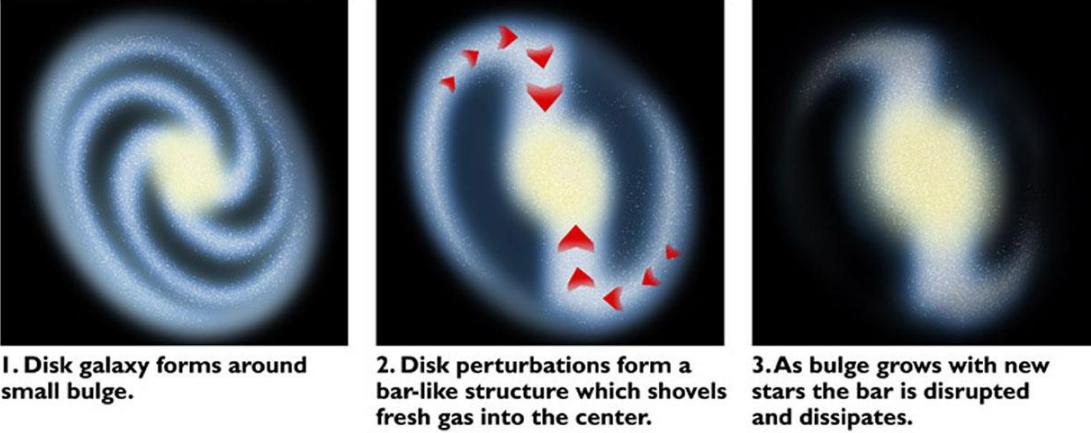
Rapid Collapse

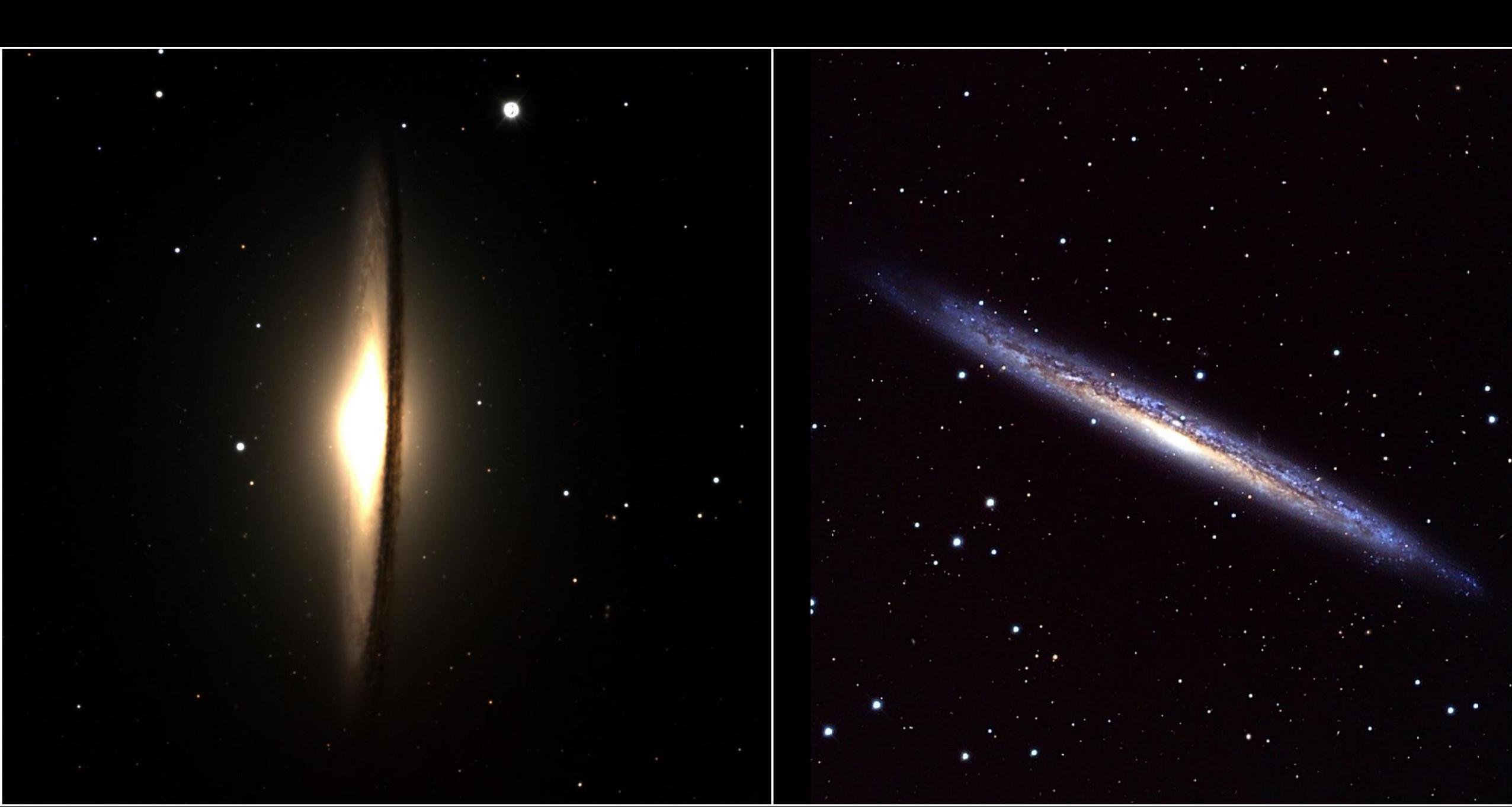


Environmental Effects

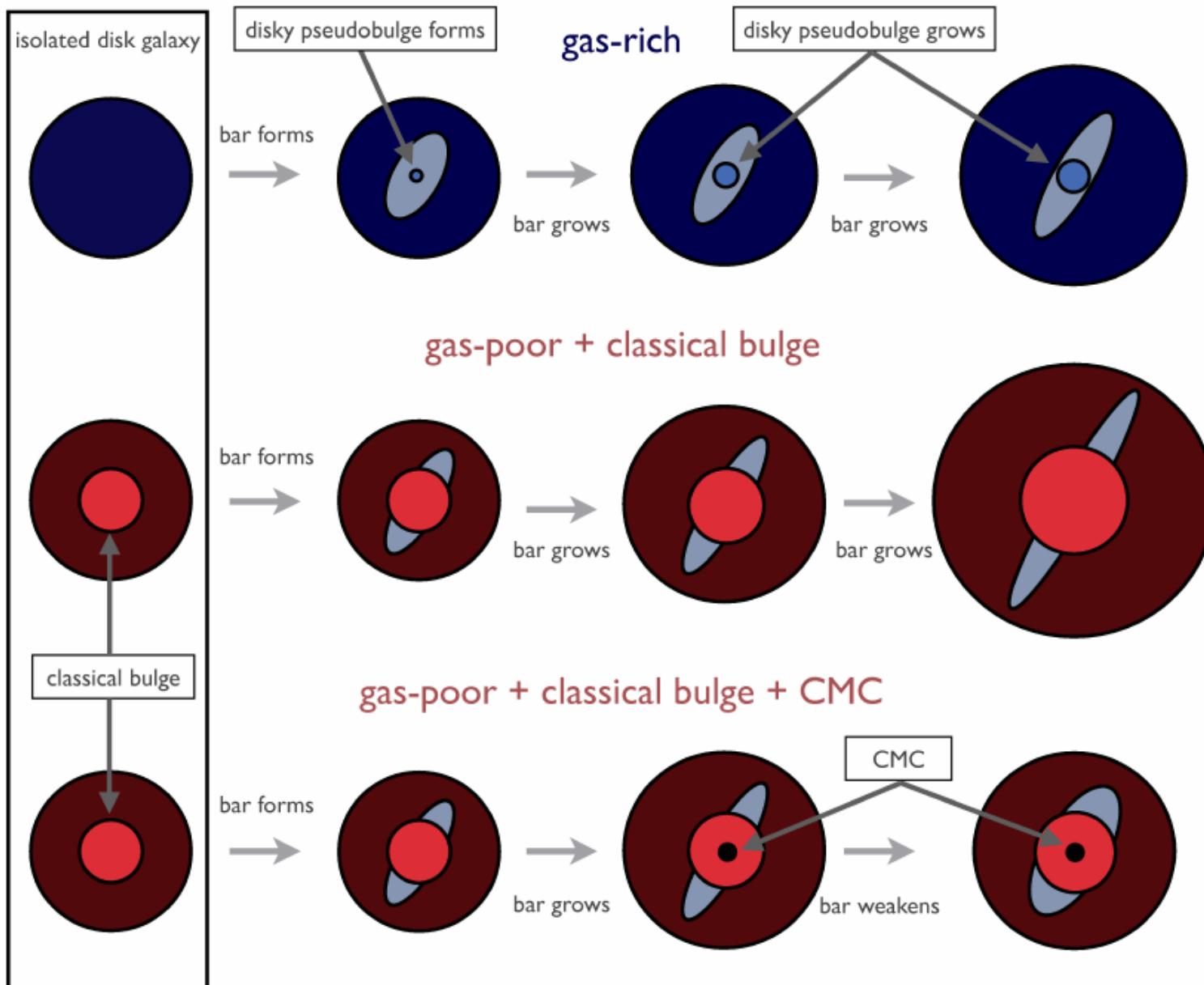


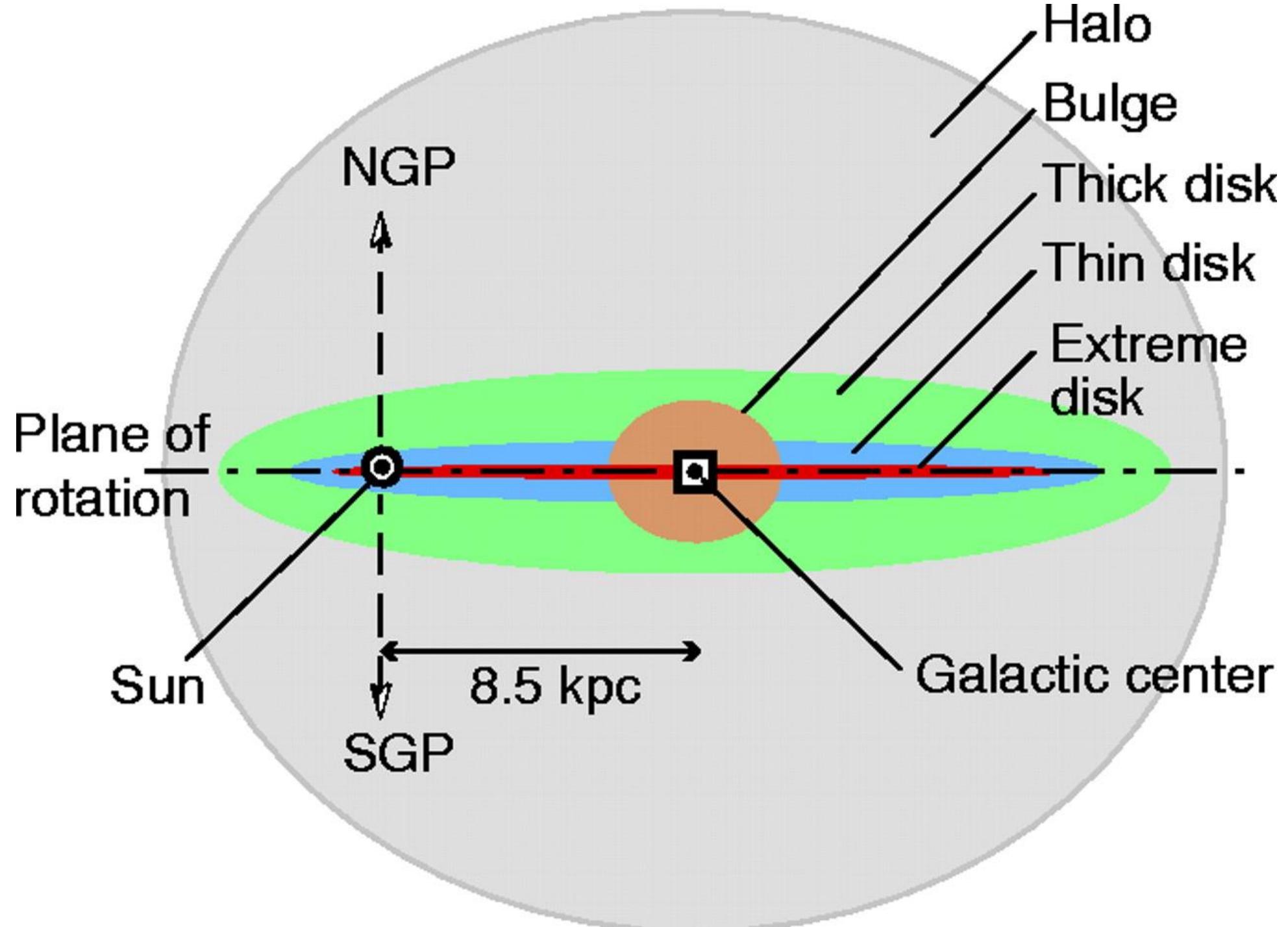
Internal Evolution

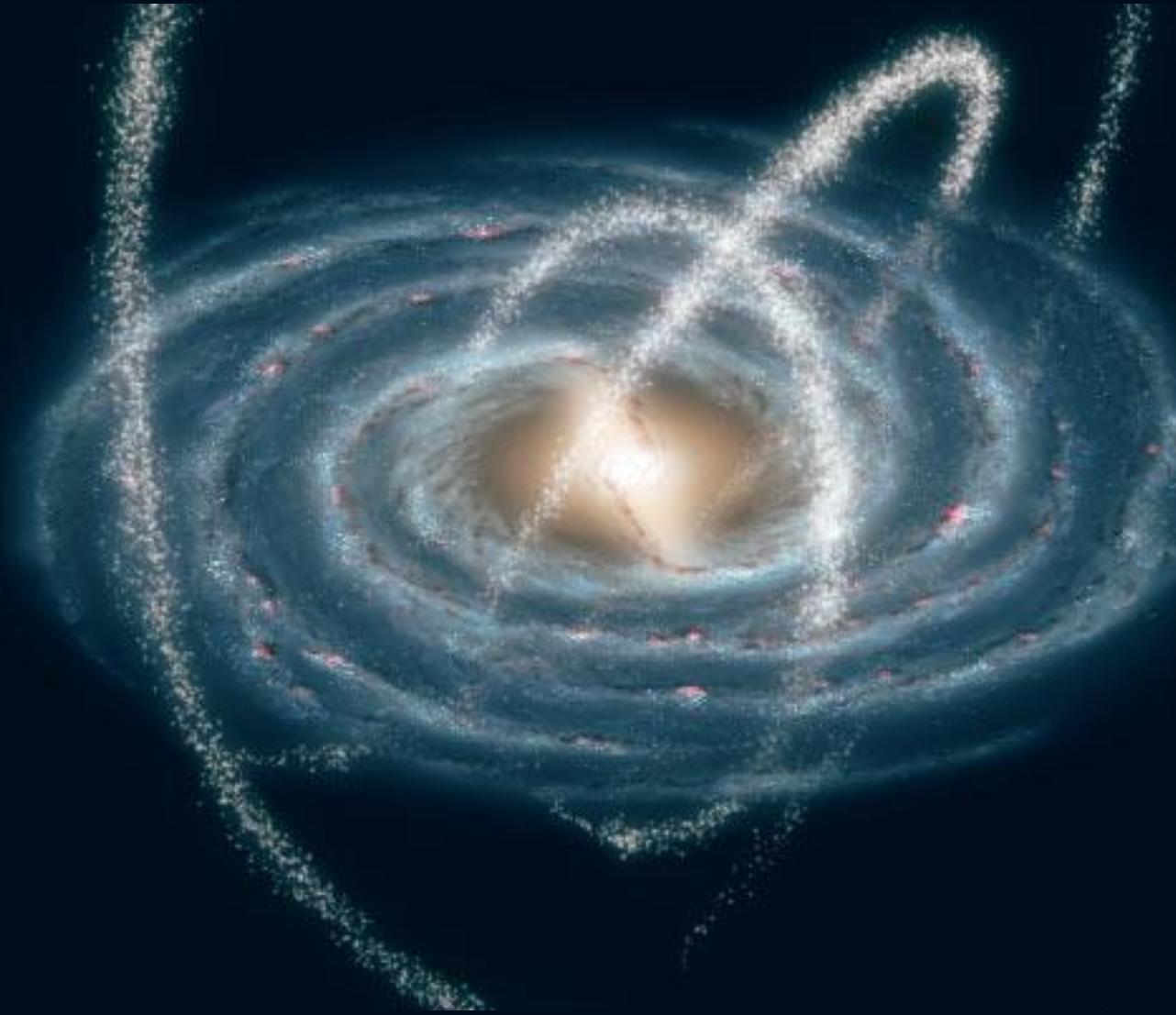


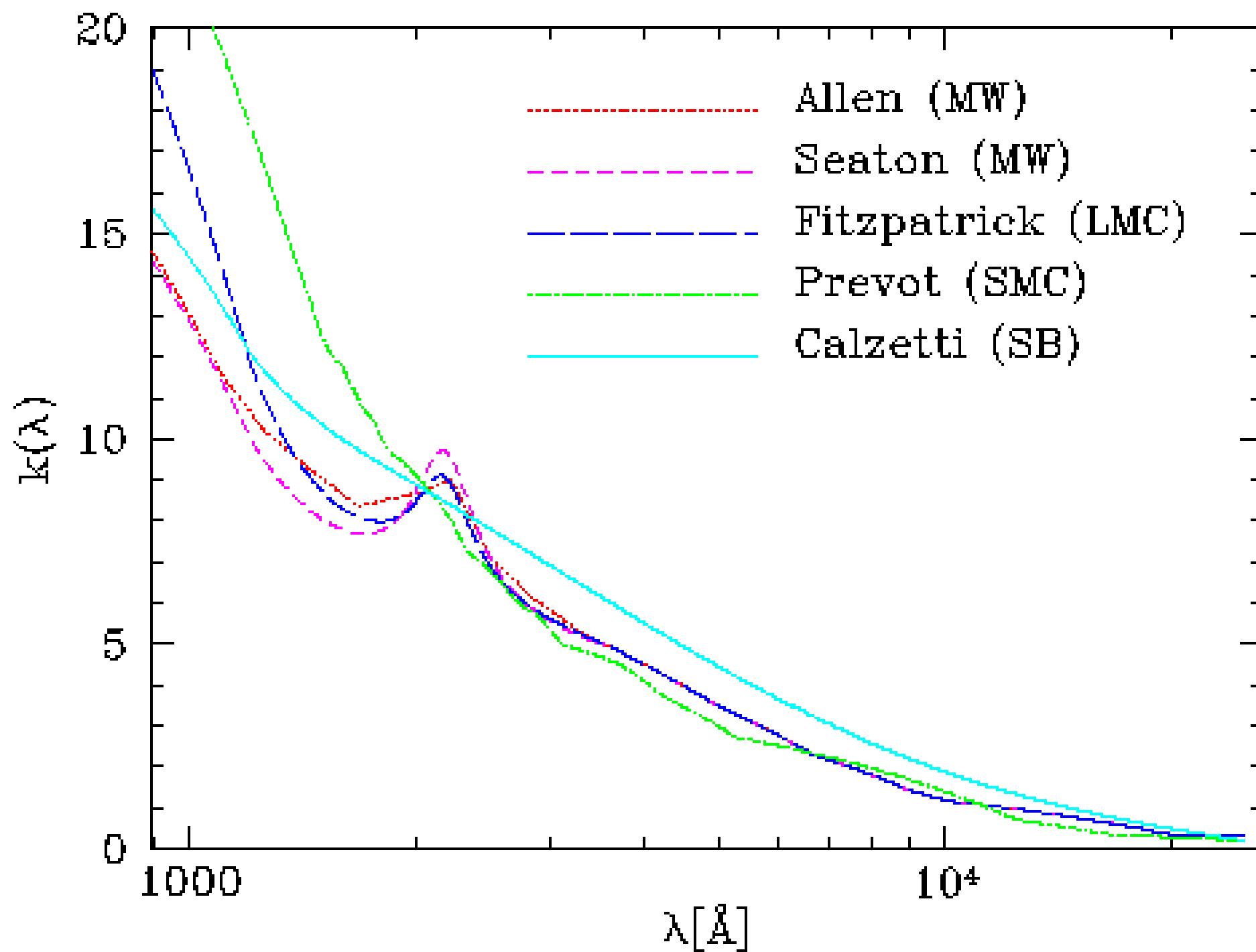


Schematic of Bar Formation and Evolution

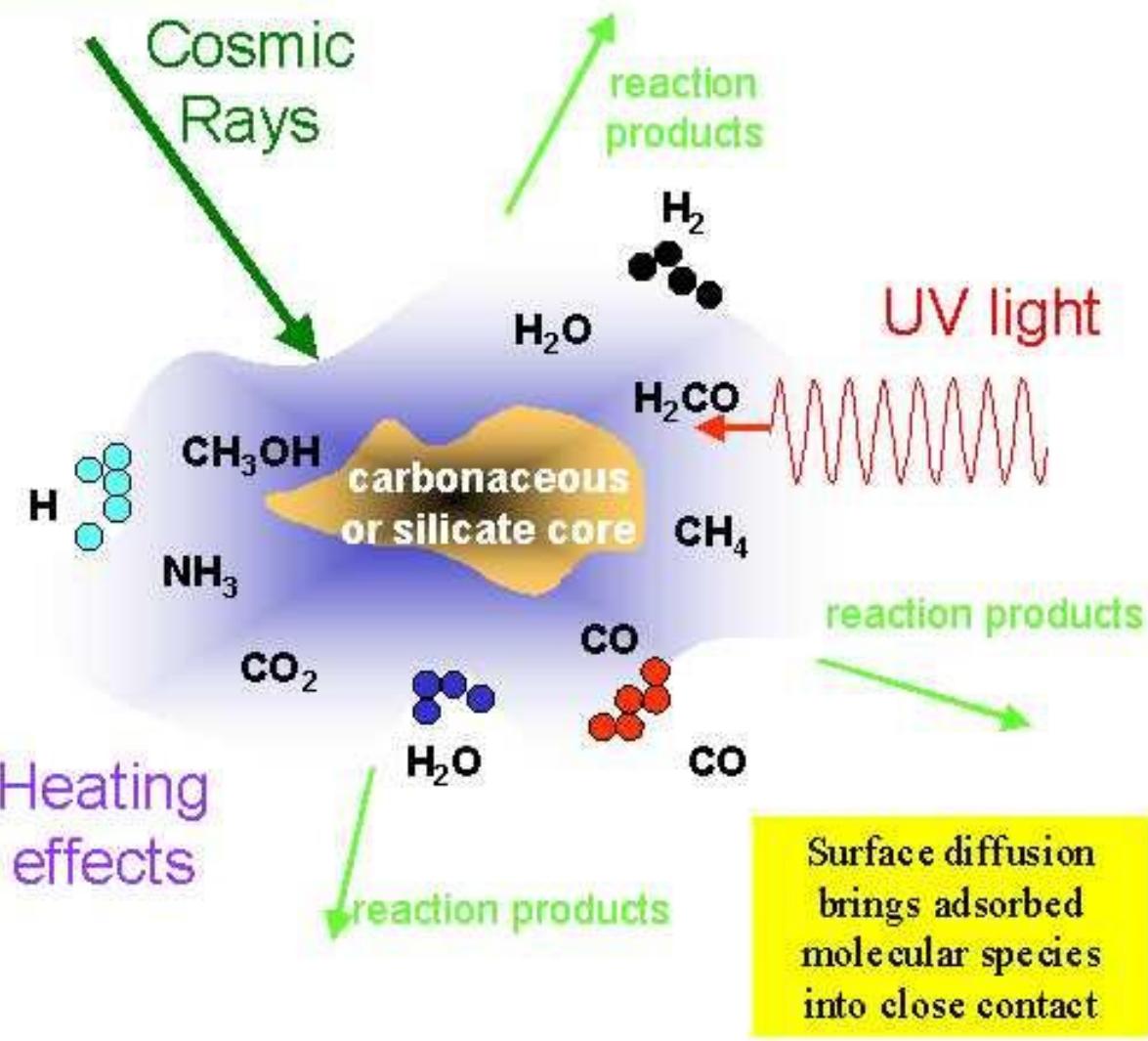


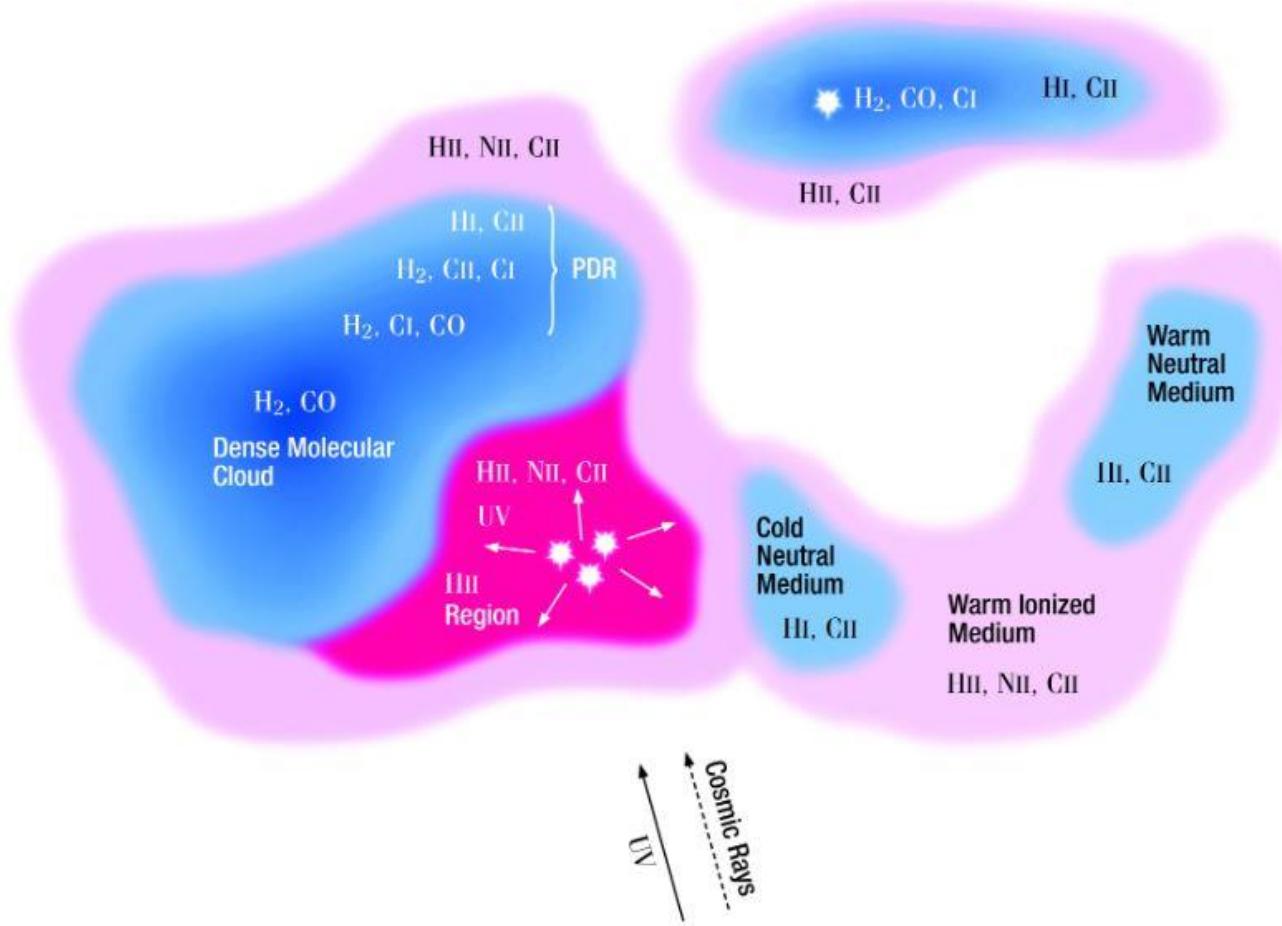






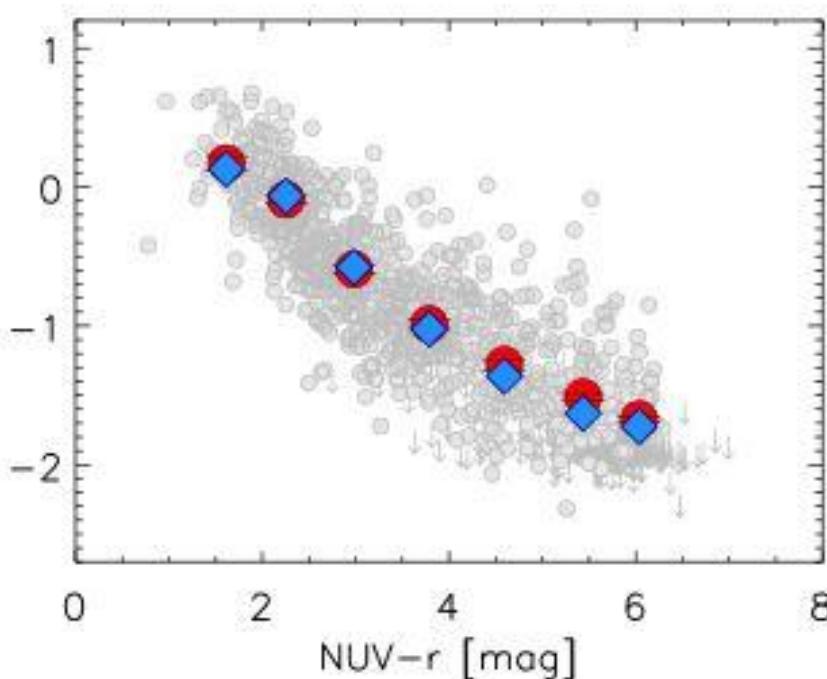
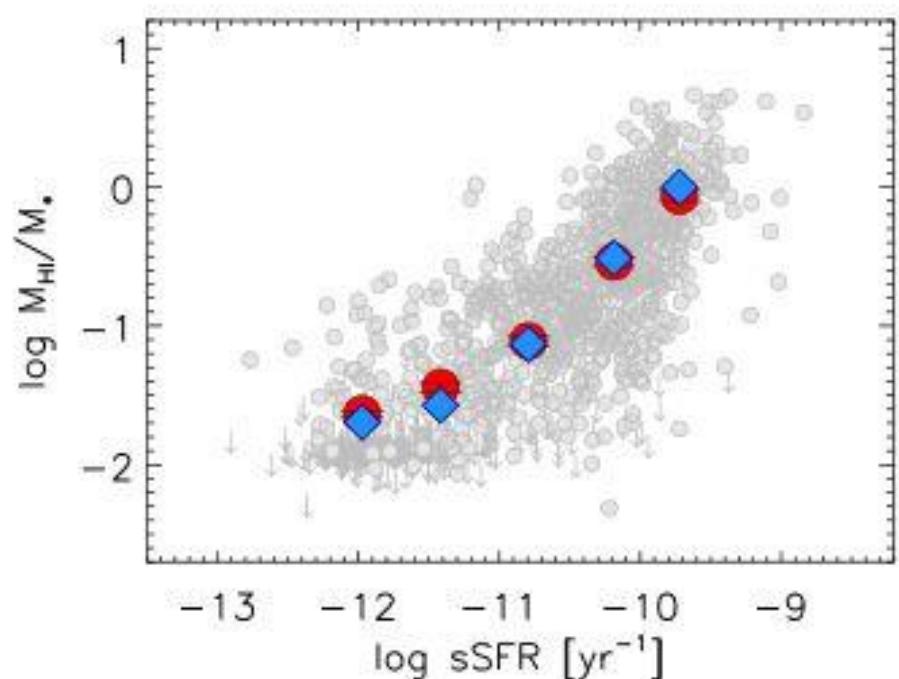
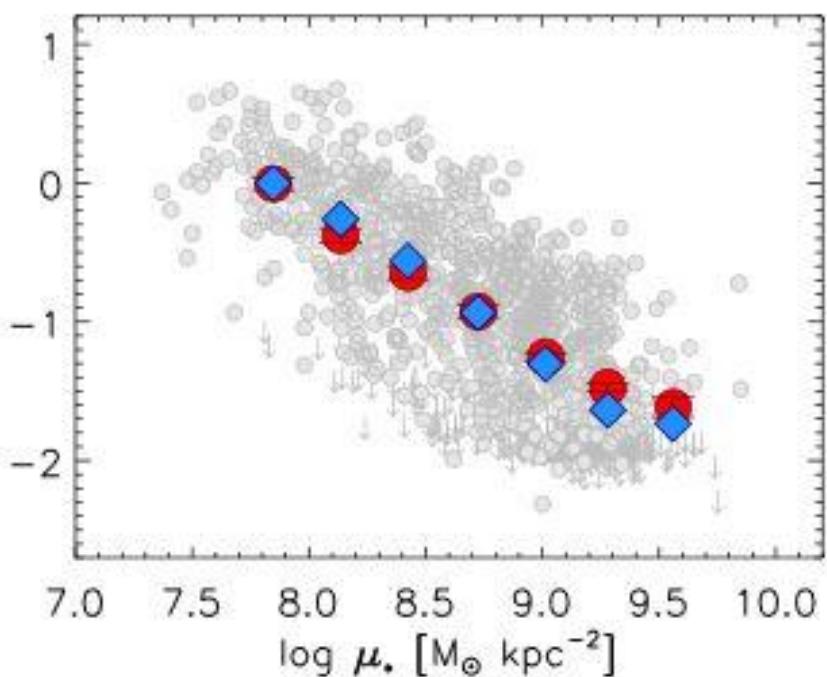
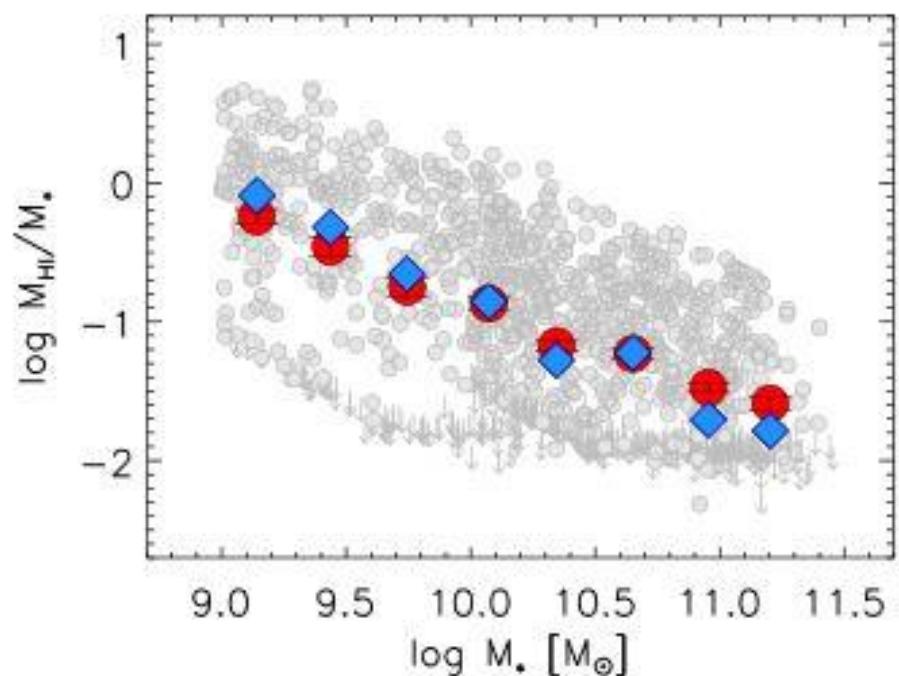
Grain provides a
'catalytic' surface
with a weak H-
bonded network

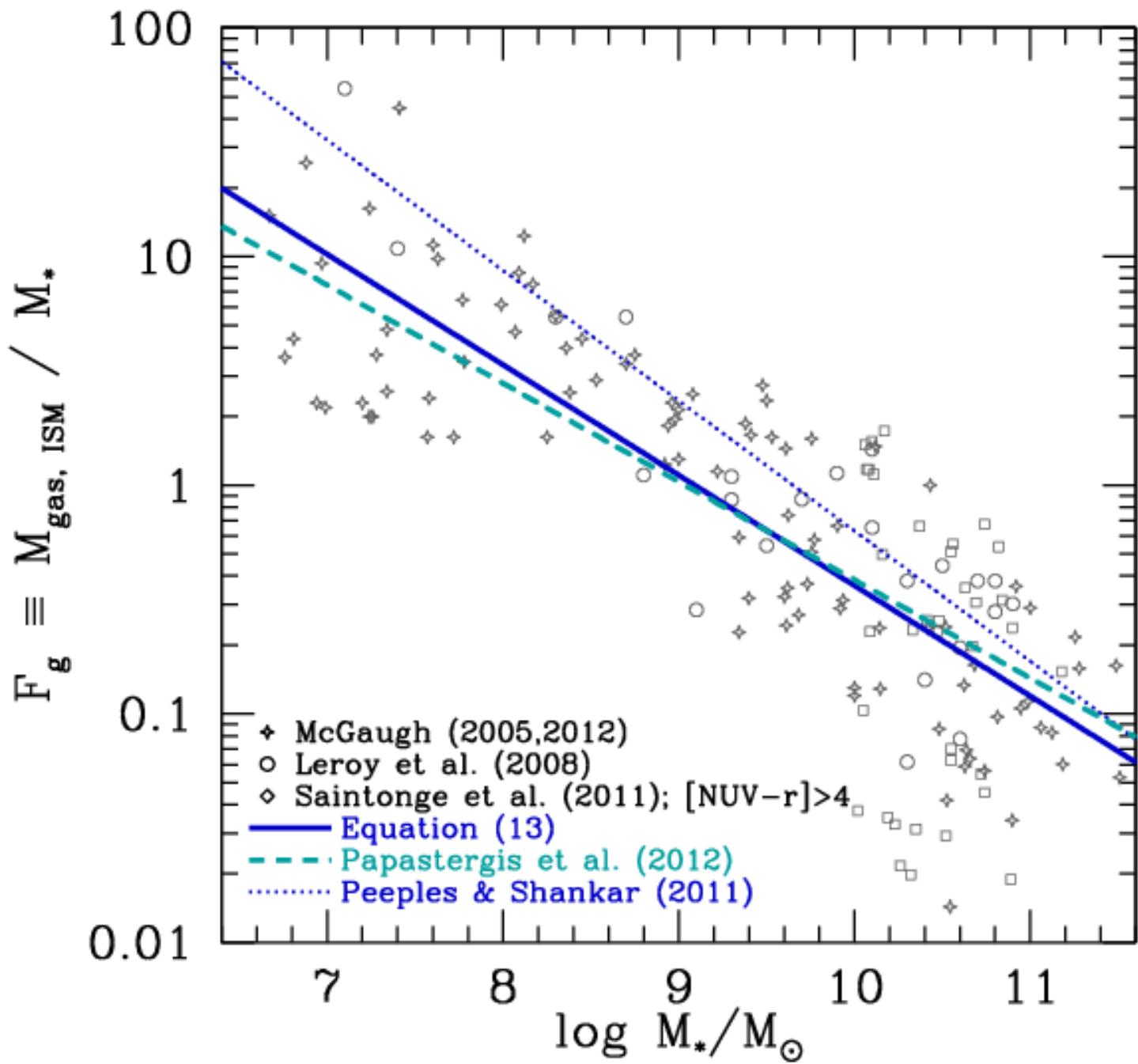


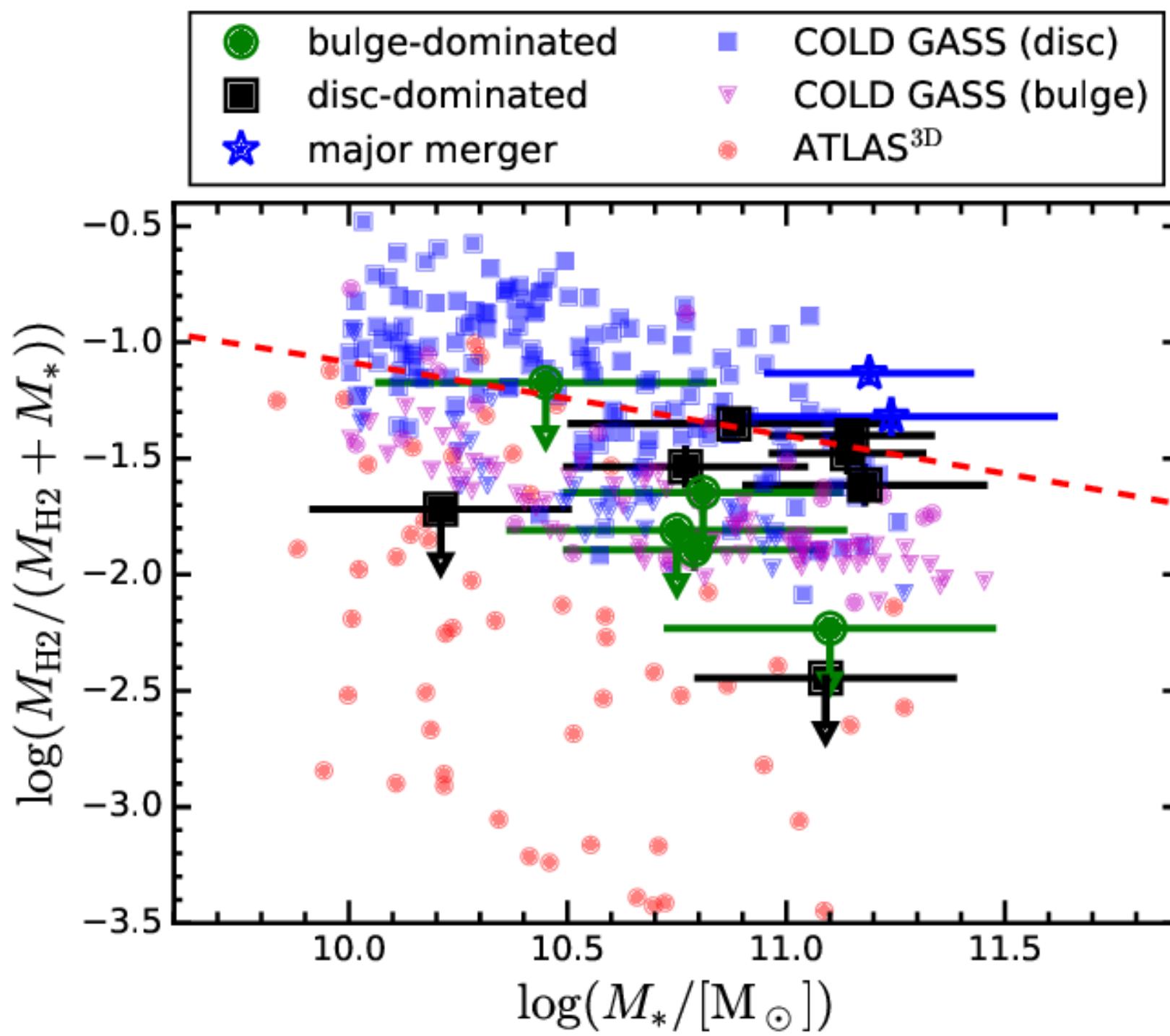


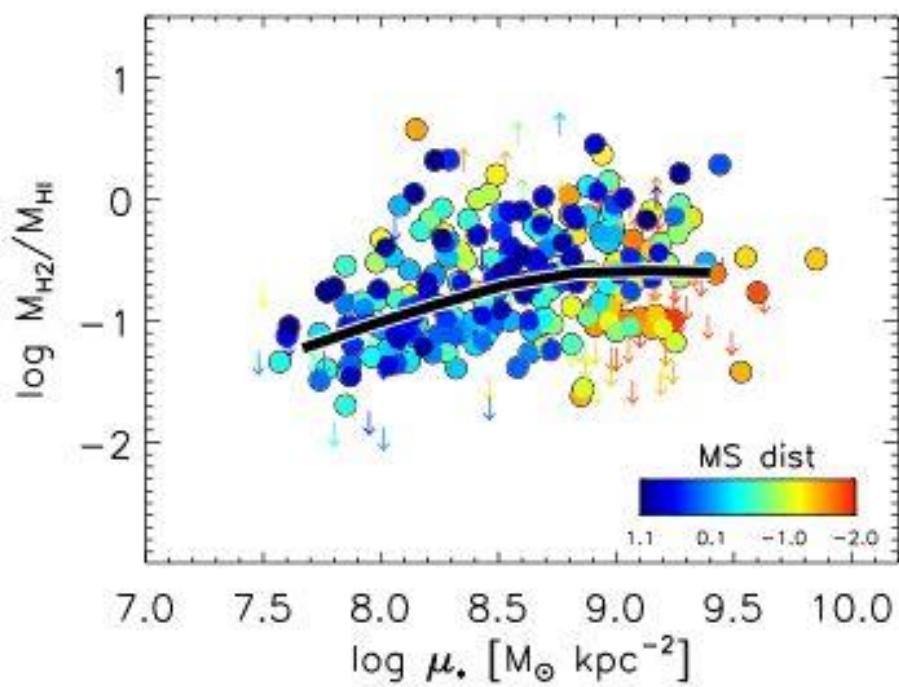
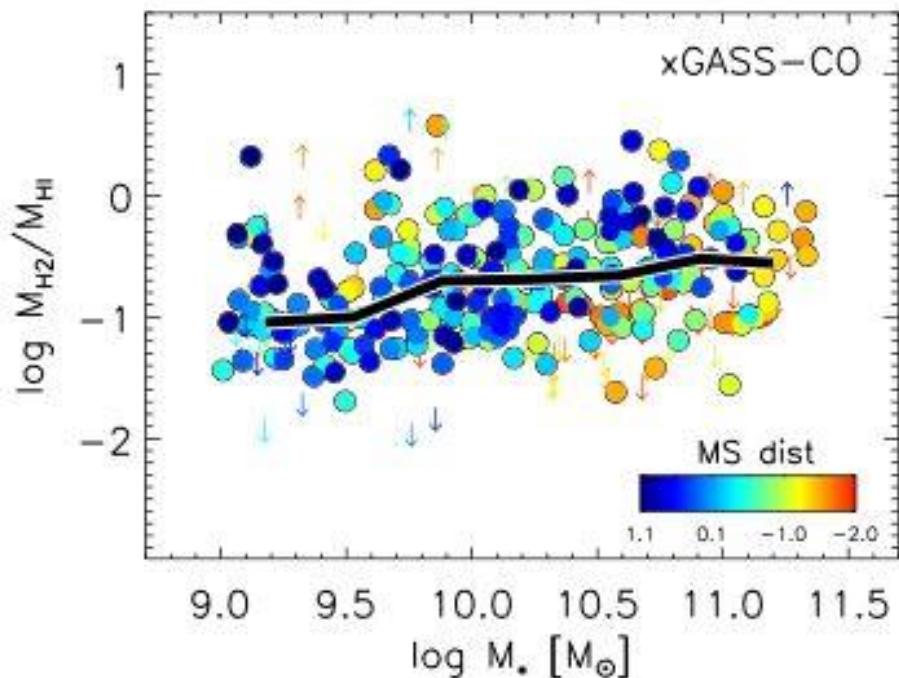
$$P(r) = \frac{\rho(r)K_B T(r)}{m}$$

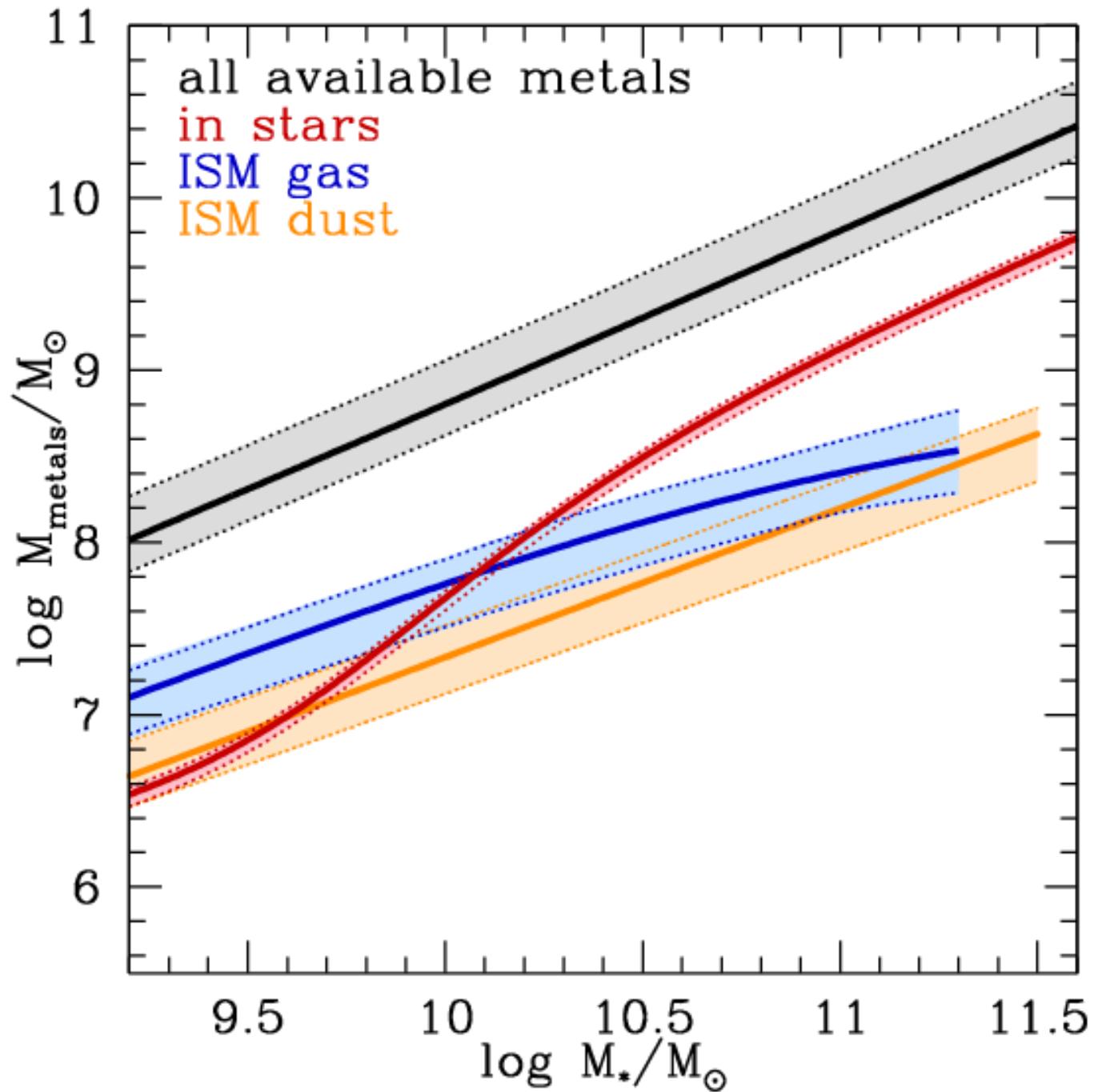
Astronomical Name	Density (cm^{-3})	Temperature (K)	β	filling factor
Molecular Cloud	$200 - \geq 10^5$	≤ 100	9×10^{-8}	0.050 %
Cold Neutral Medium (CNM)	10 - 100	~ 100	...	1 %
HII regions	5 - 10	8000	15 - 30	3×10^{-3} %
Warm Neutral Medium (WNM)	0.1 - 0.5	8000	...	30 %
Warm Ionized Medium (WIM)	0.1 - 0.5	8000	0.29	25 %
Very Local Medium (VLISM)	0.11	6700	0.27	6 - 19 %
Coronal (HIM)	5×10^{-3}	10^6	1.8	50 %

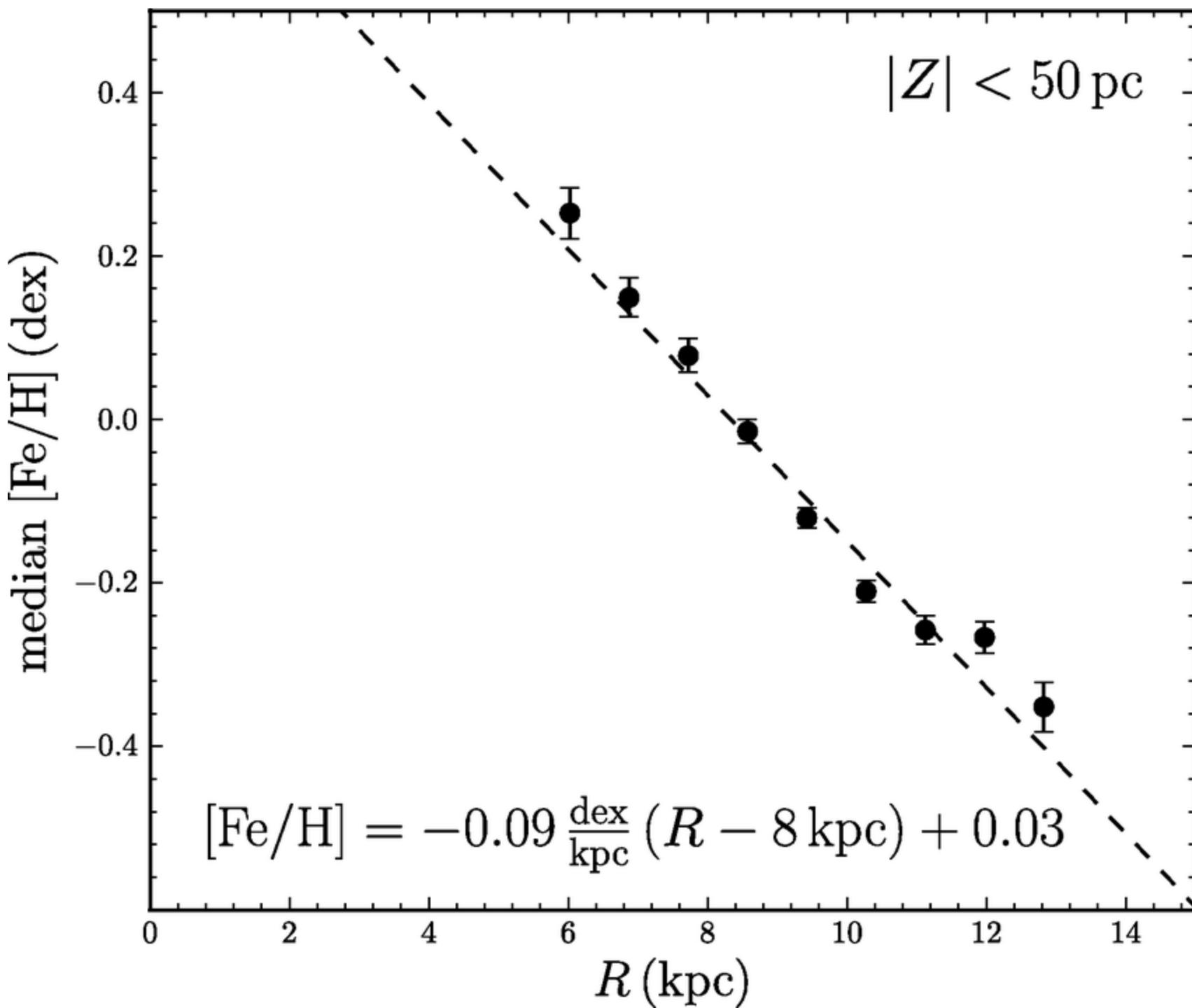


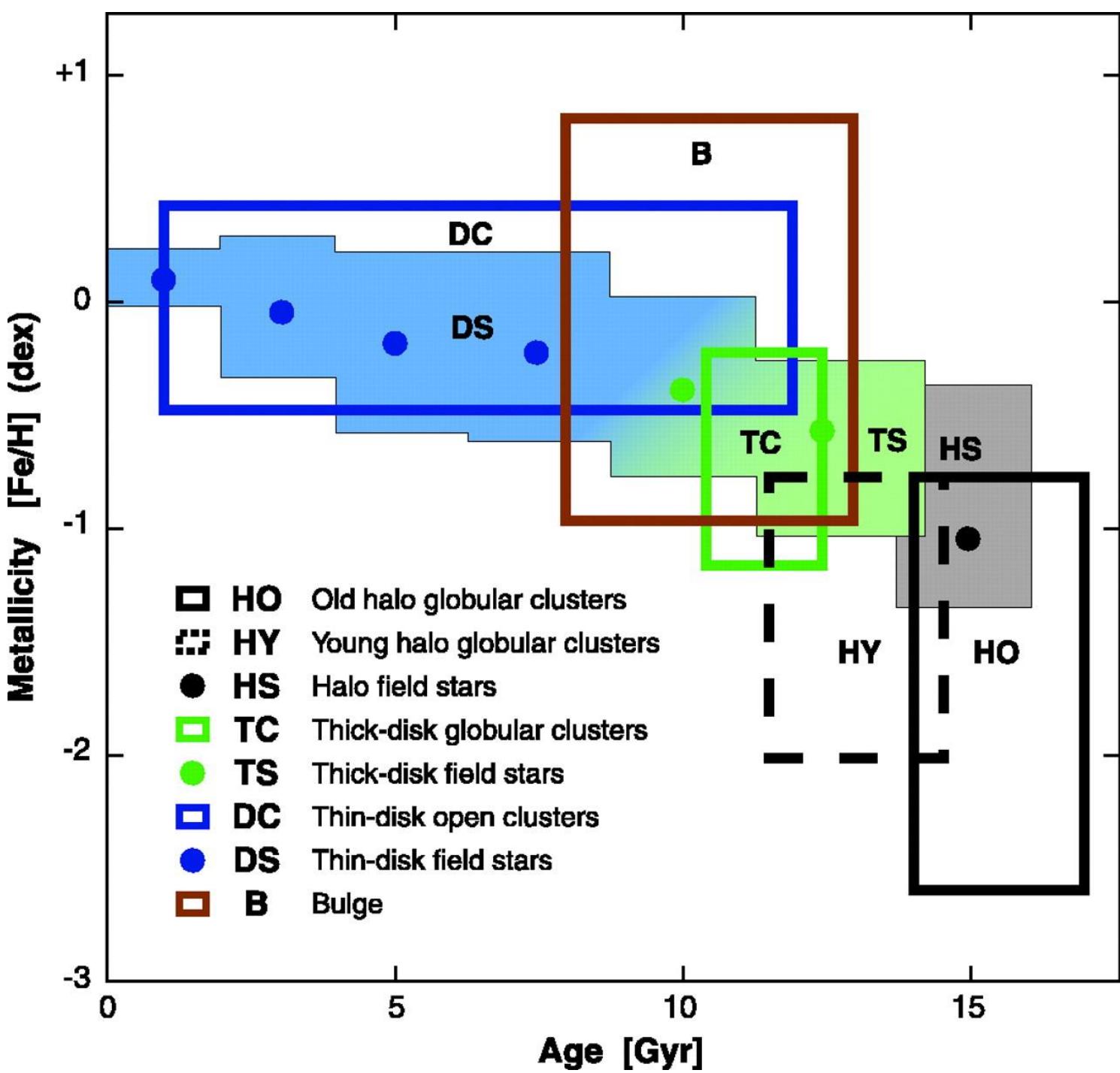




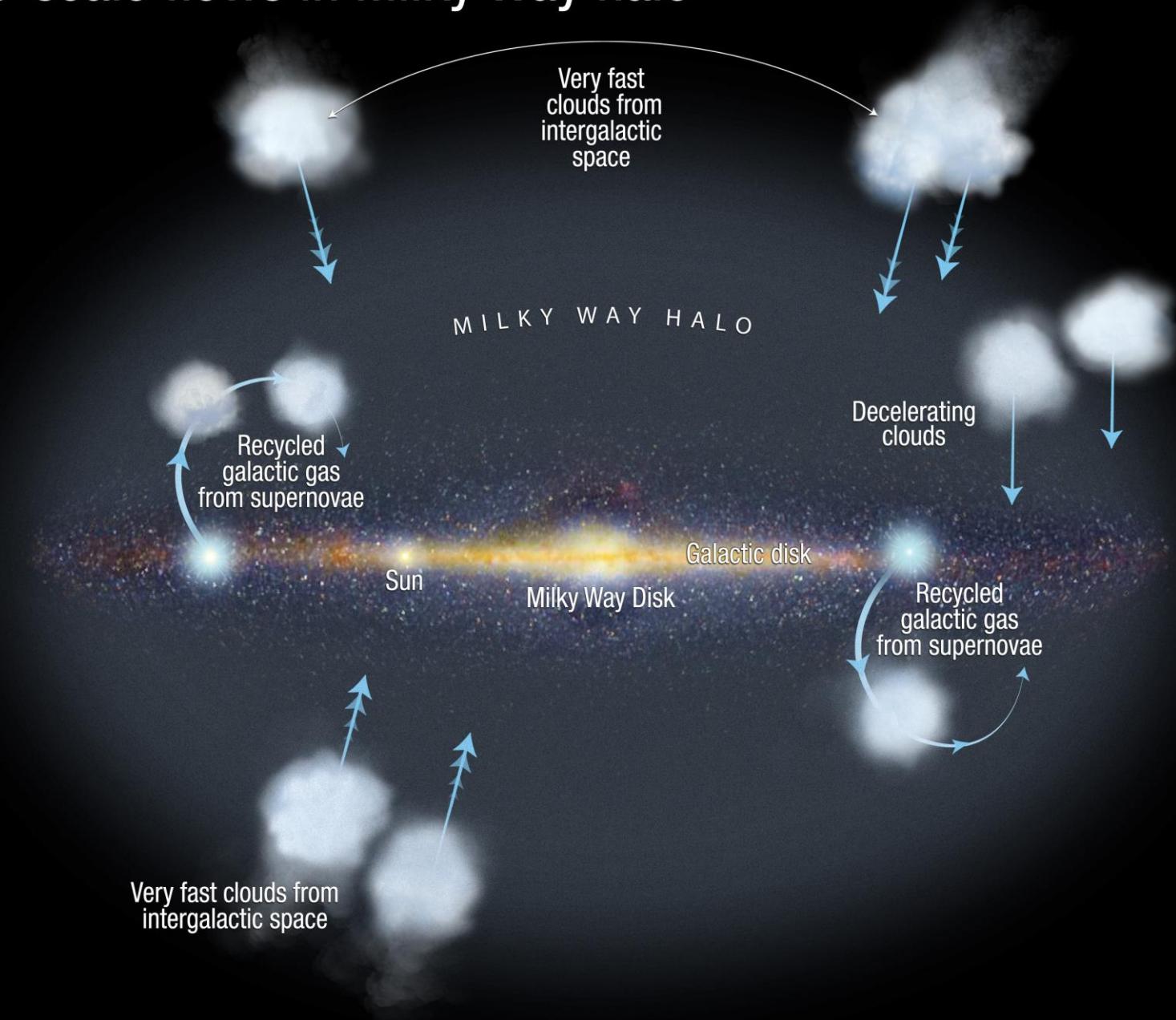


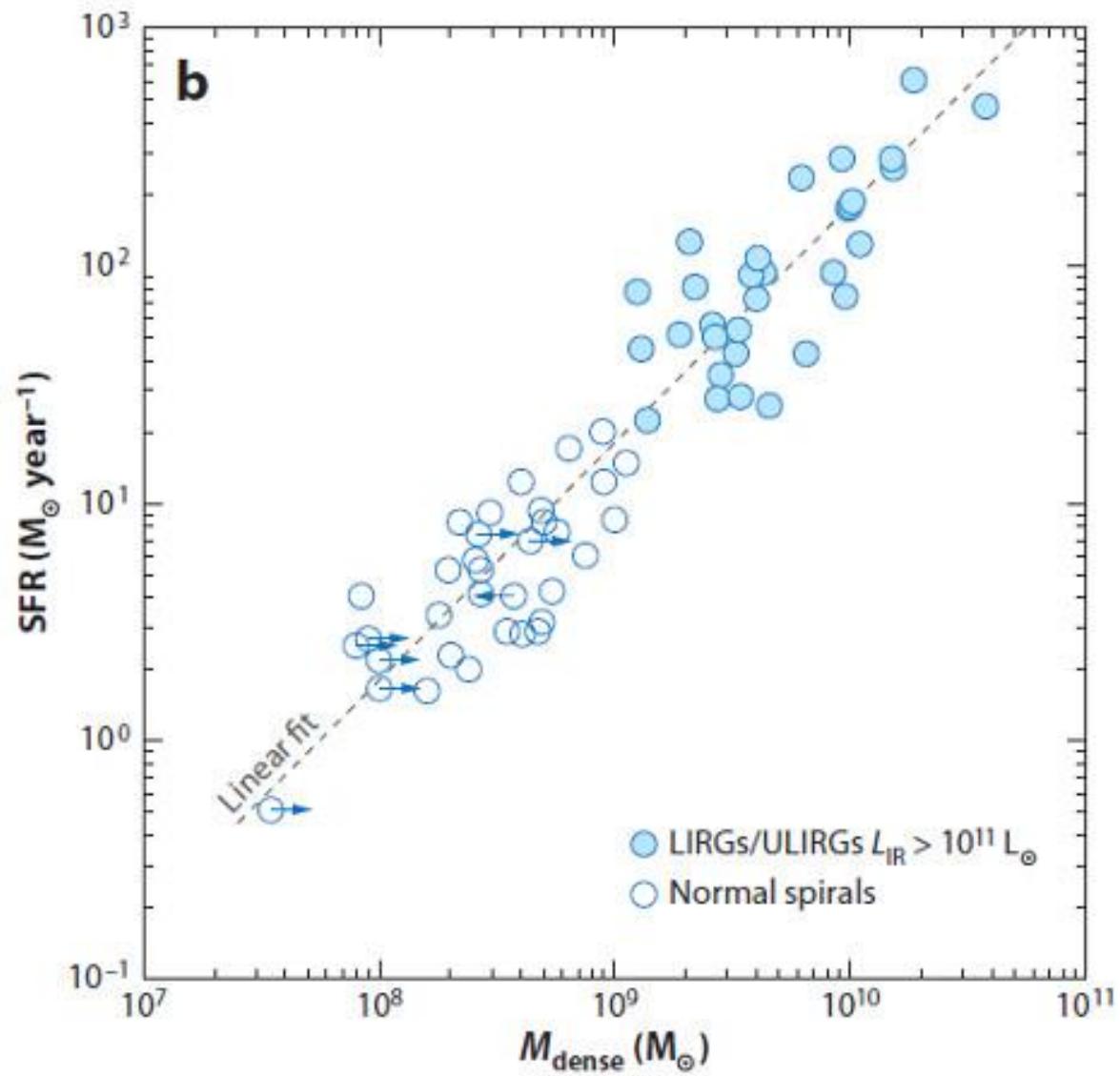
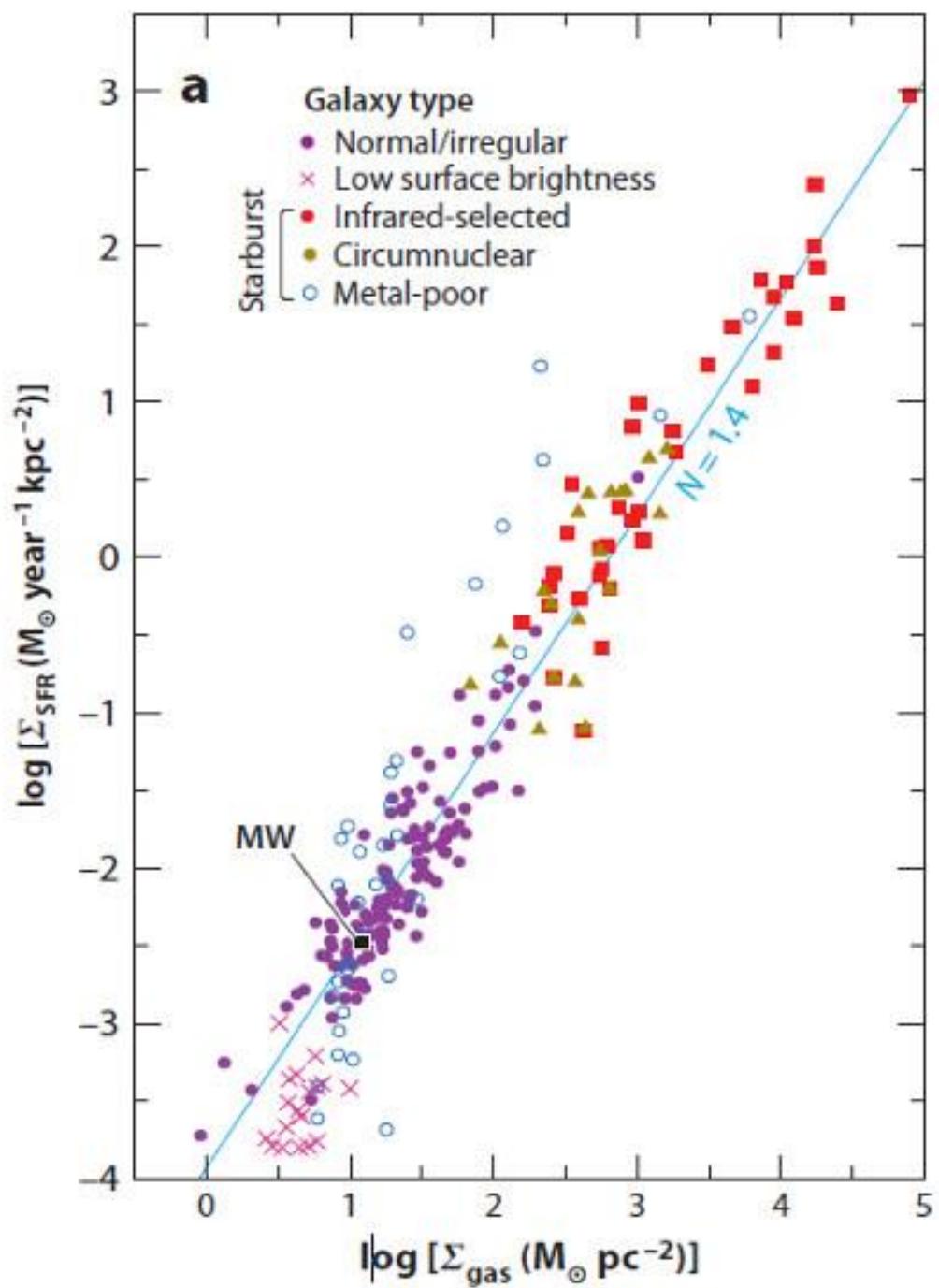


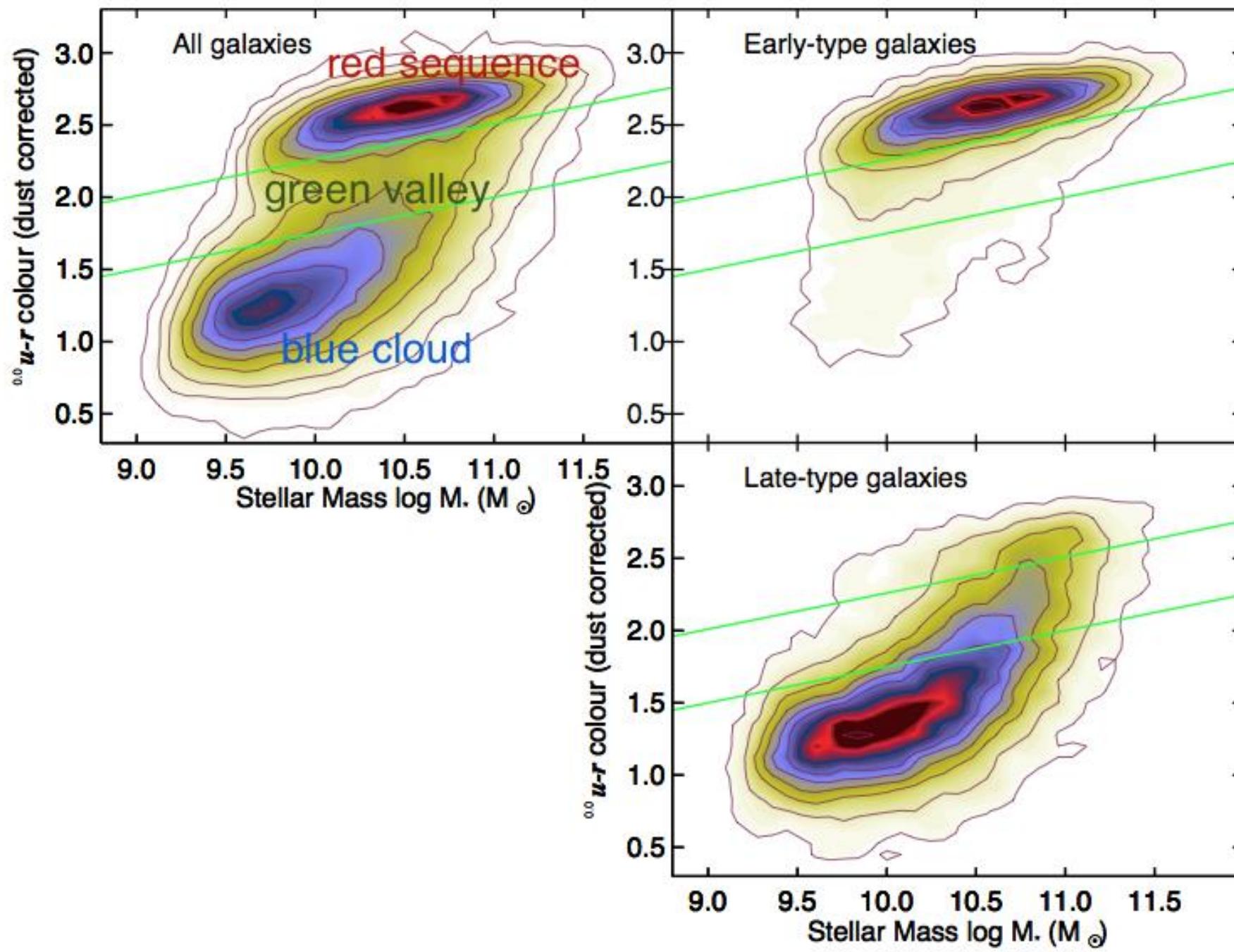


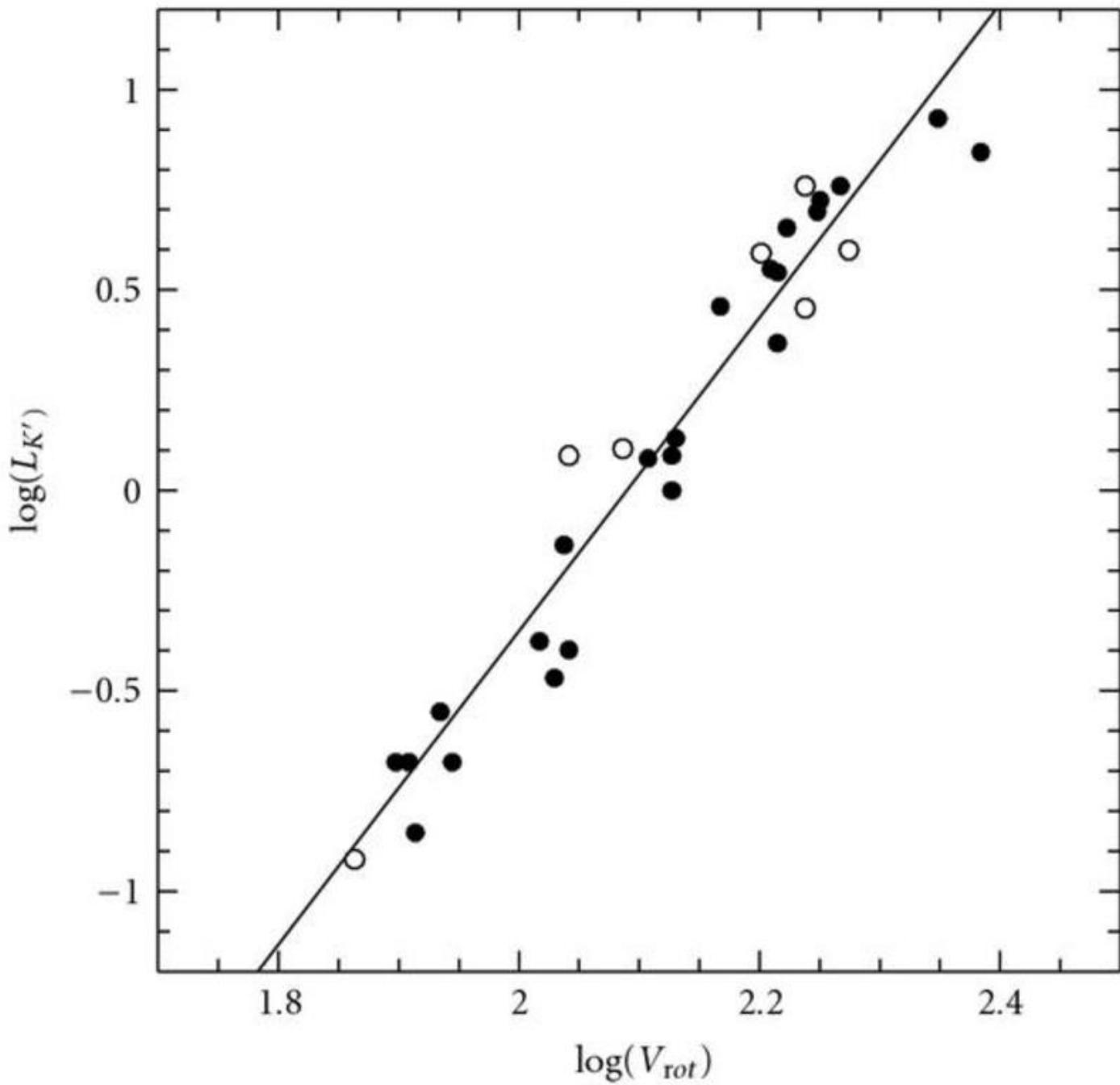


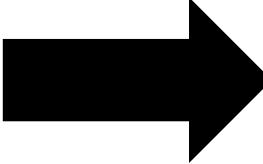
Large-scale flows in Milky Way halo







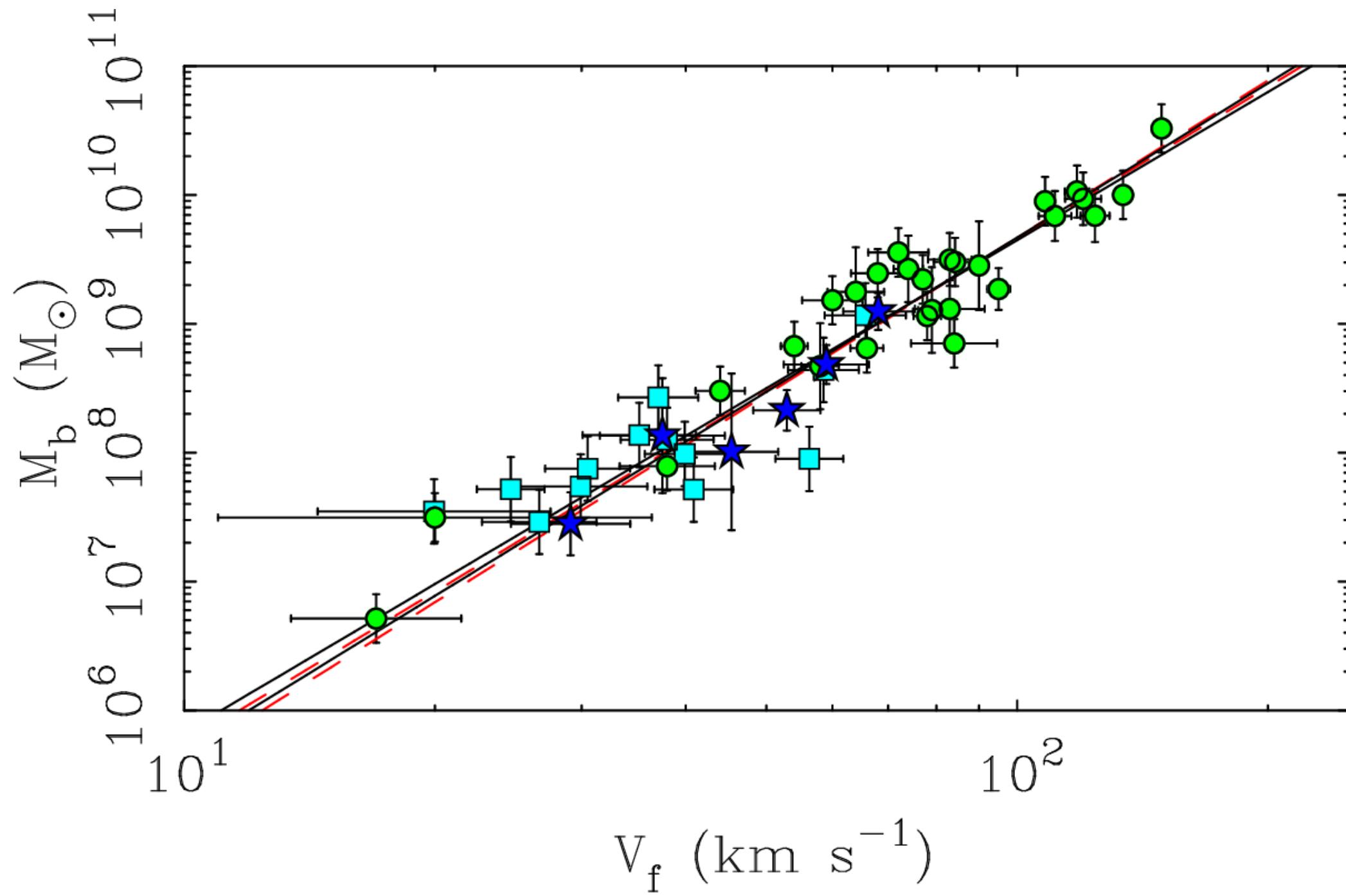


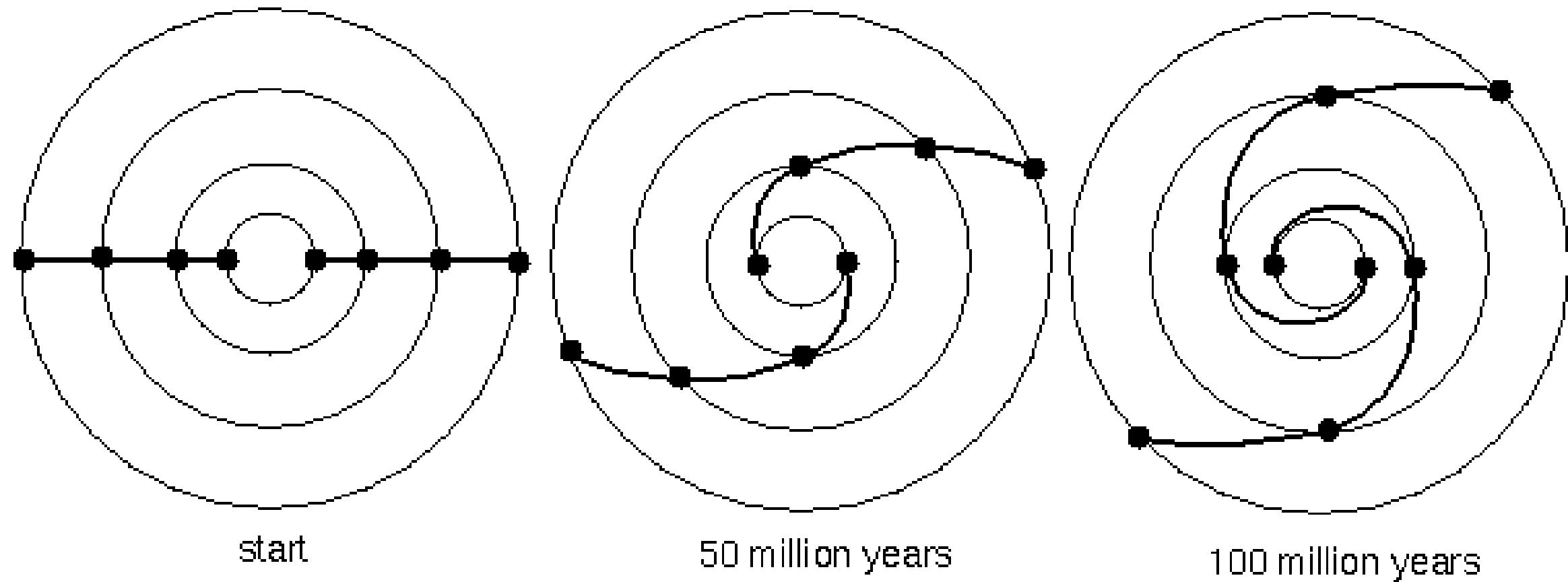
Assunzioni: $\begin{cases} \frac{M}{L} = \frac{M_*}{L} \left(1 + \frac{M_{\text{oscura}}}{M_*} \right) \approx \text{costante} \\ \mu_0(\text{banda B}) \approx 21.7 \frac{\text{mag}}{\text{arcsec}^2} \end{cases}$ 

$L = k v_{\max}^4$

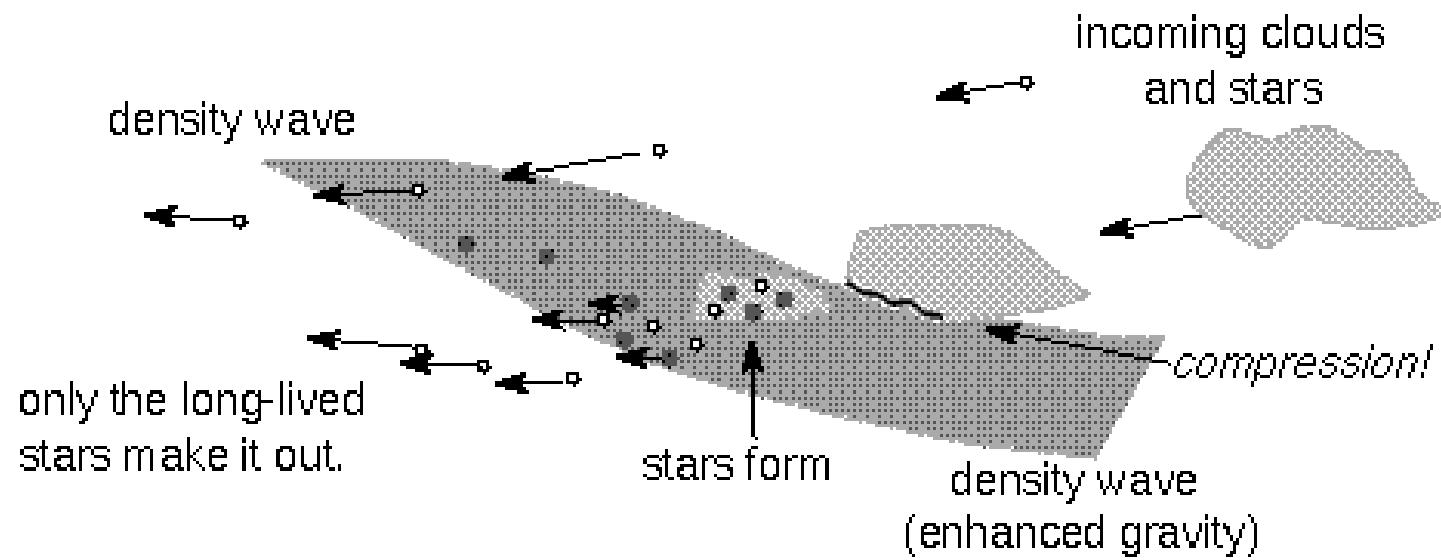
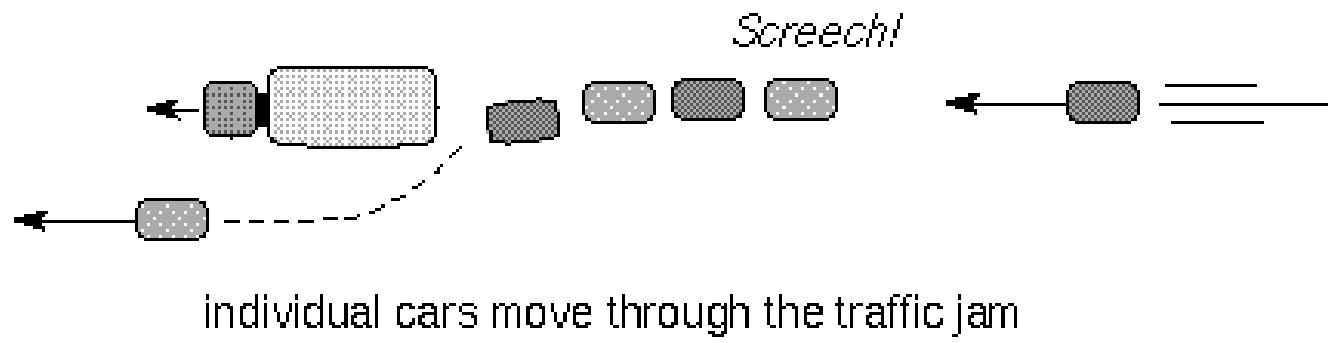
$$\begin{cases} v_{\max}^2 \propto \frac{M}{R}, \text{ si bilanciano forza centrifuga e gravitazionale} \\ \mu_0 \propto \frac{L}{R^2}, \text{ per definizione} \end{cases}$$

$$(v_{\max}^2)^2 \propto \left(\frac{M}{R}\right)^2 \propto \frac{M^2}{L} \mu_0 \propto \left(\frac{M}{L}\right)^2 \mu_0 L \propto L$$

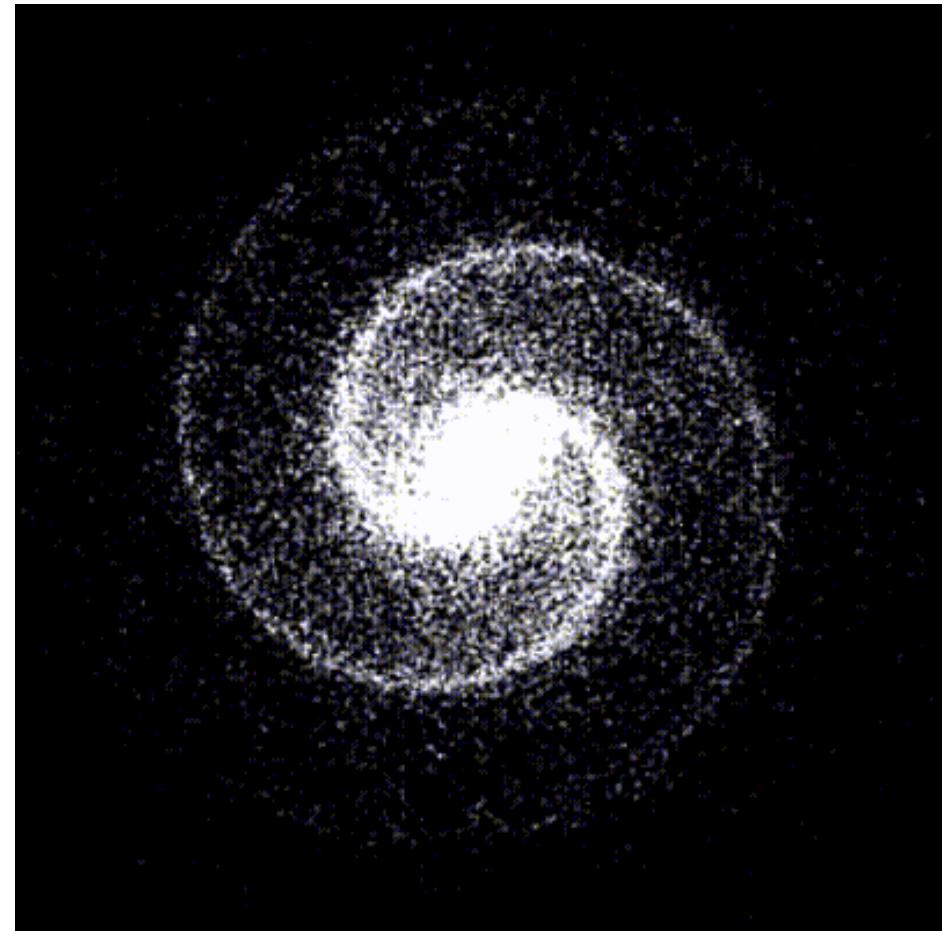




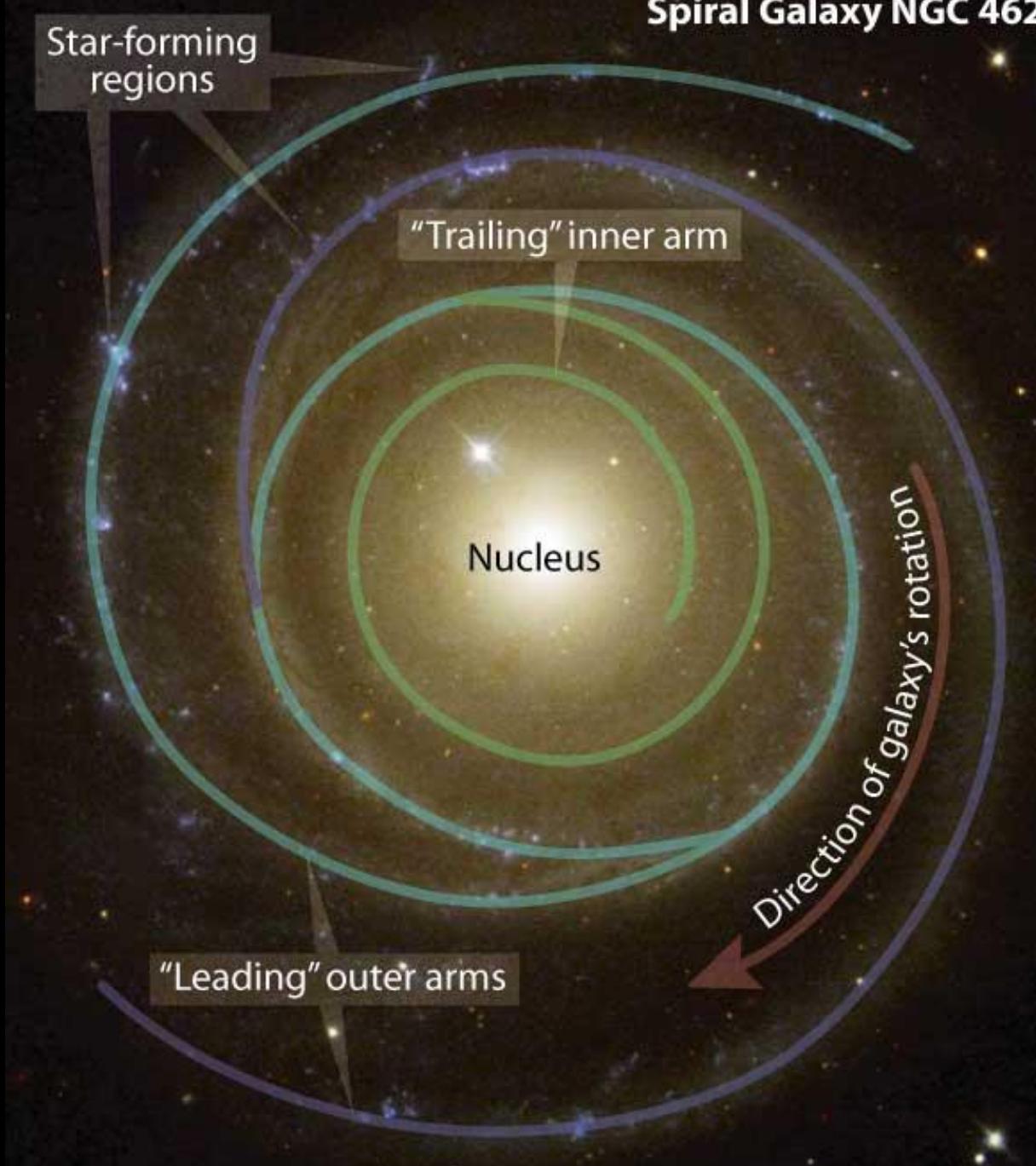
Differential rotation: stars near the center take less time to orbit the center than those farther from the center. Differential rotation can create a spiral pattern in the disk in a short time.

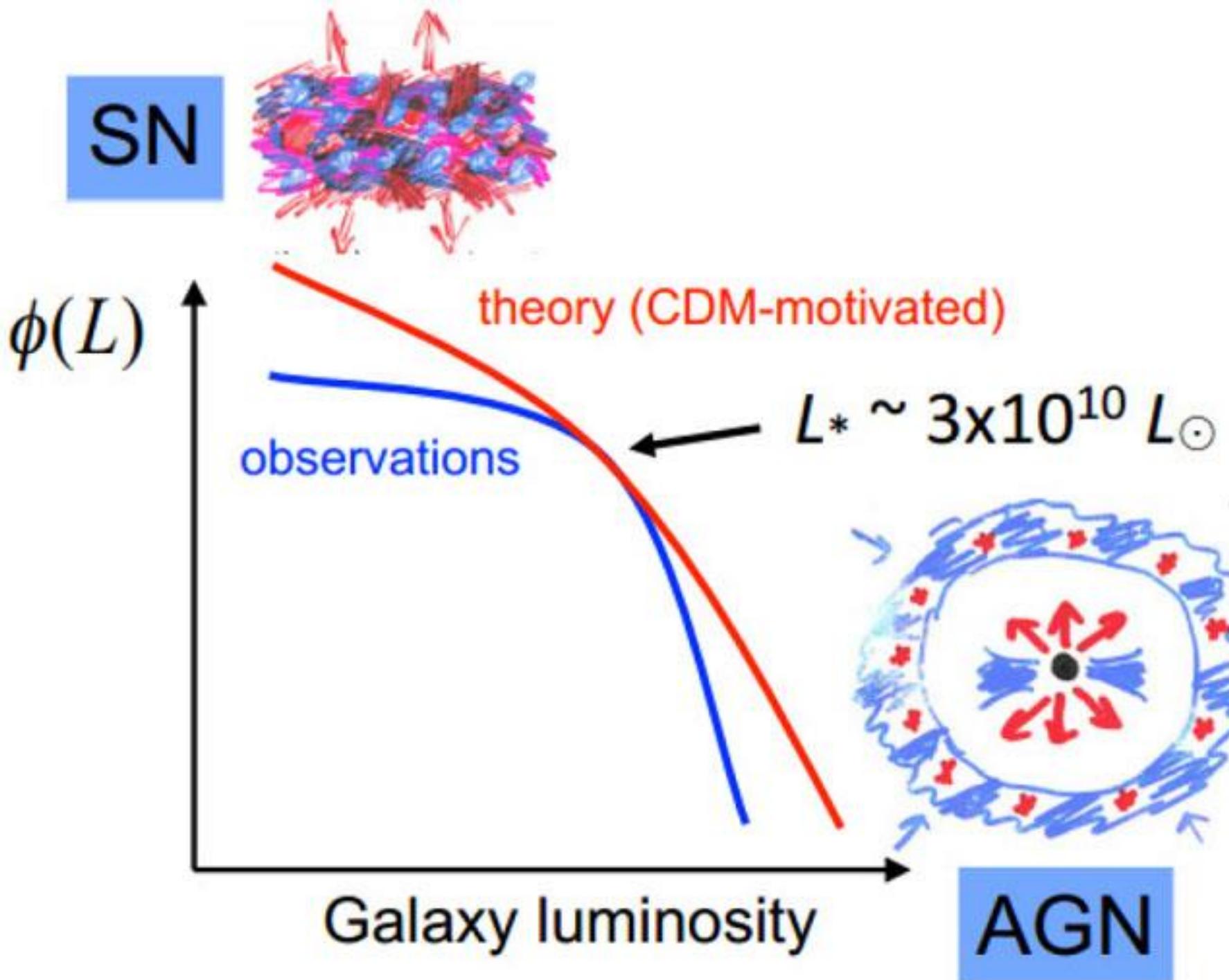


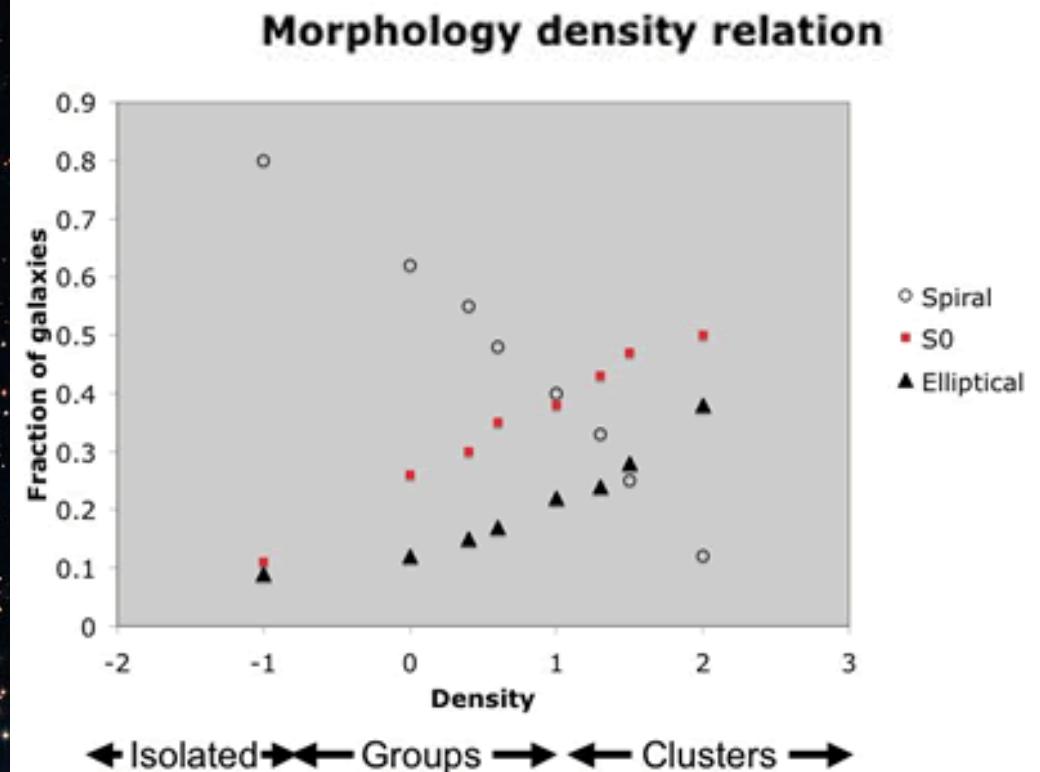
Spiral density waves are like traffic jams. Clouds and stars speed up to the density wave (are accelerated toward it) and are tugged backward as they leave, so they accumulate in the density wave (like cars bunching up behind a slower-moving vehicle). Clouds compress and form stars in the density wave, but only the fainter stars live long enough to make it out of the wave.



Spiral Galaxy NGC 4622









Dinamica nelle ETG

Examples:

	N	R	v	τ_{cross}	τ_{relax}	age/ τ_{relax}
open cluster	100	2 pc	0.5 km/s	$4 \cdot 10^6$ yrs	10^7 yrs	≥ 1
globular cluster	10^5	4 pc	10 km/s	$4 \cdot 10^5$ yrs	$4 \cdot 10^8$ yrs	≥ 10
ellipt. galaxy	10^{12}	10 kpc	600 km/s	$2 \cdot 10^7$ yrs	10^{17} yrs	10^{-7}
dwarf galaxy	10^9	1 kpc	50 km/s	$2 \cdot 10^7$ yrs	10^{14} yrs	10^{-4}
galaxy cluster	1000	1 Mpc	1000 km/s	10^9 yrs	$2 \cdot 10^{10}$ yrs	10^{-1}



Conclusion: stellar encounters will perturb a star only over $\tau_{\text{relax}} \sim 10^{17}$ yr, or $5 \times 10^9 \tau_{\text{cross}}$

Elliptical galaxies are

collisionless system: the constituent particles move under the influence of a mean gravitational field generated by a smooth mass distribution, rather than a collection of mass points

Credit:
Pellegrini Silvia

Fast e slow rotators

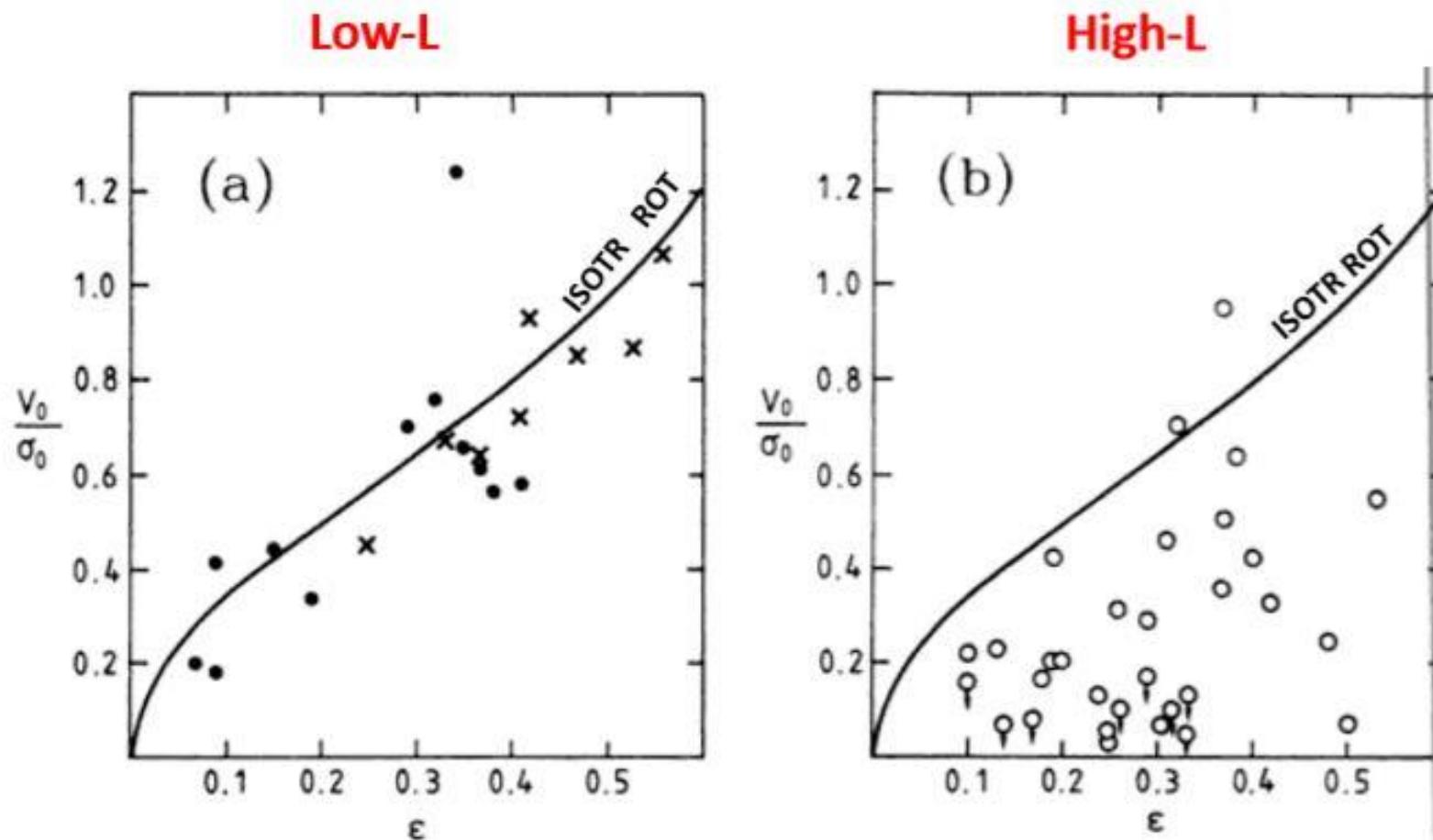
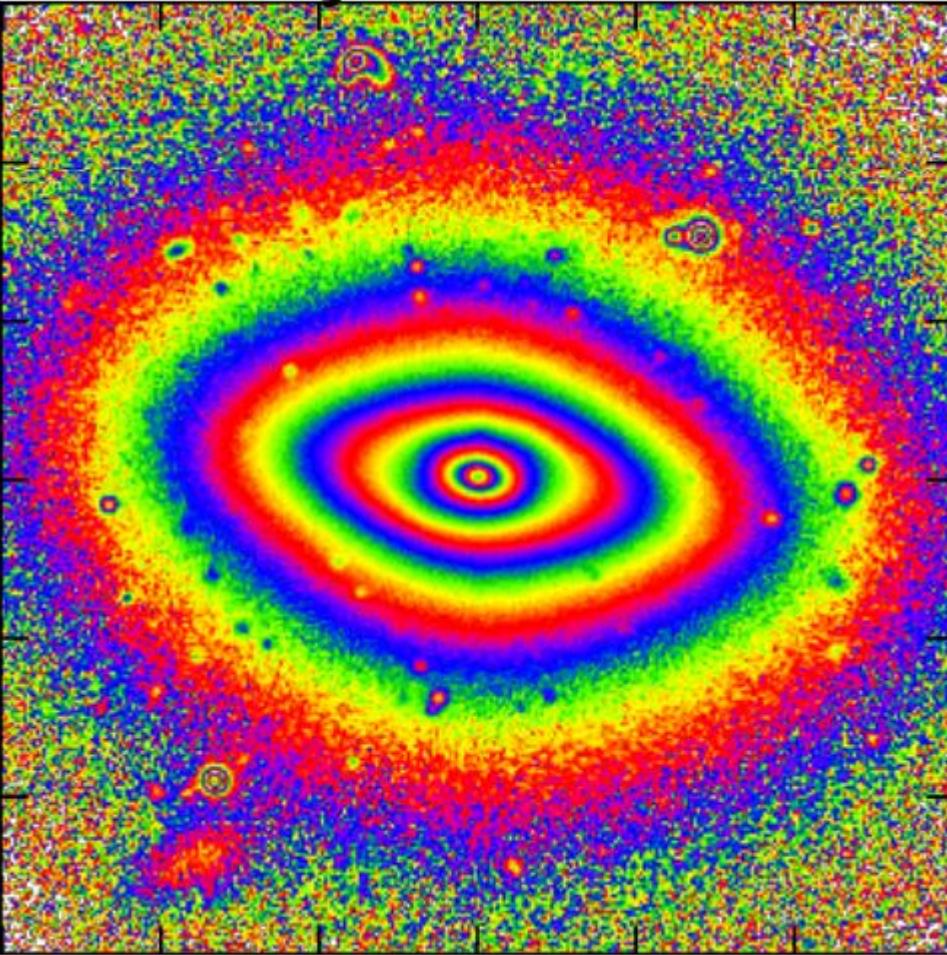


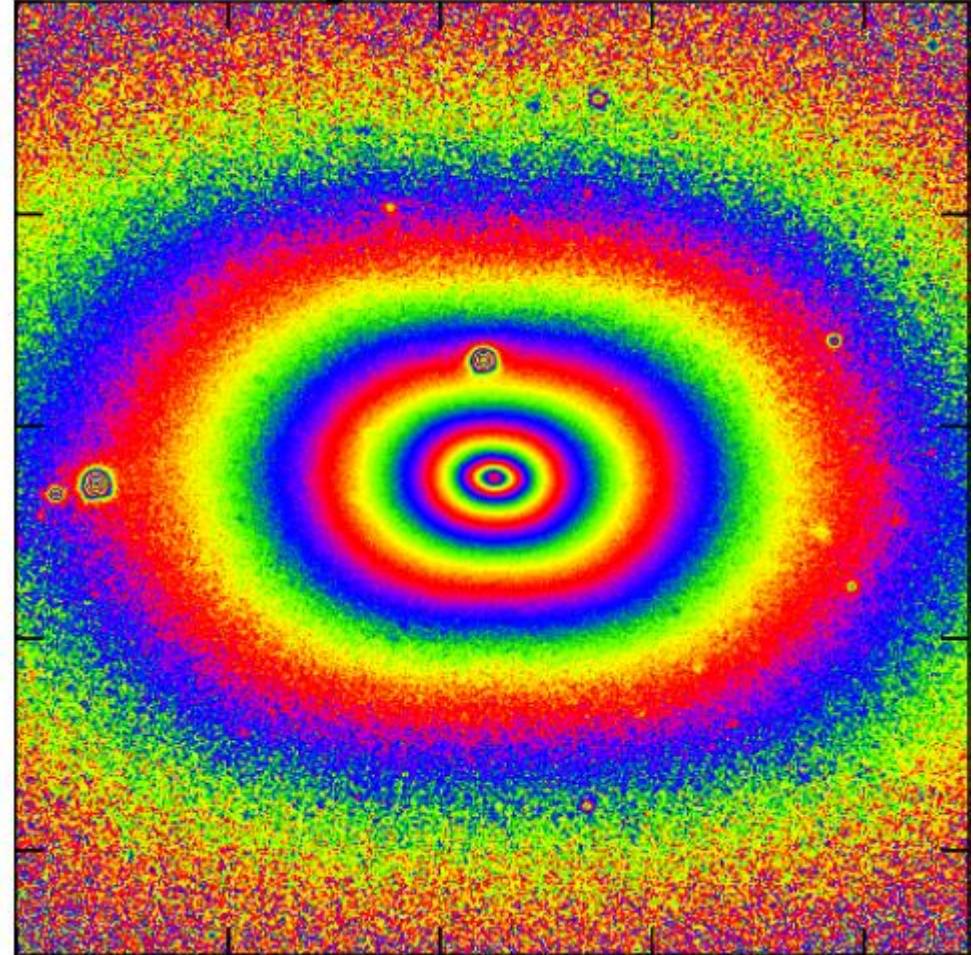
Figure 4-6. (a) The positions in the $(v/\sigma, \epsilon)$ plane of elliptical galaxies (dots), and of spheroids (crosses), that have luminosities smaller than $L = 2.5 \times 10^{10} L_\odot$. (b) The same as (a) but for elliptical galaxies brighter than $L = 2.5 \times 10^{10} L_\odot$. (After Davies et al. 1983.)

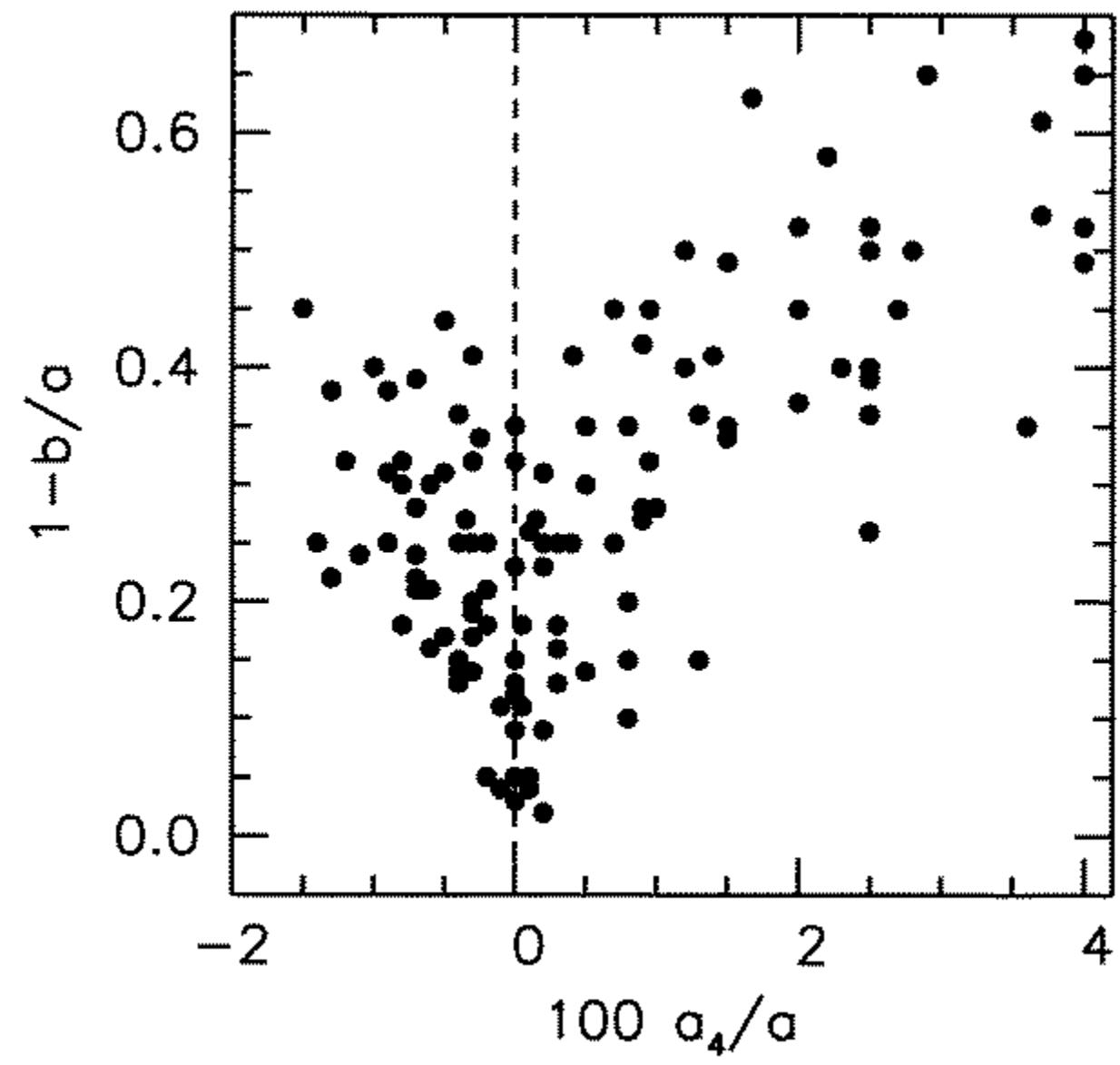
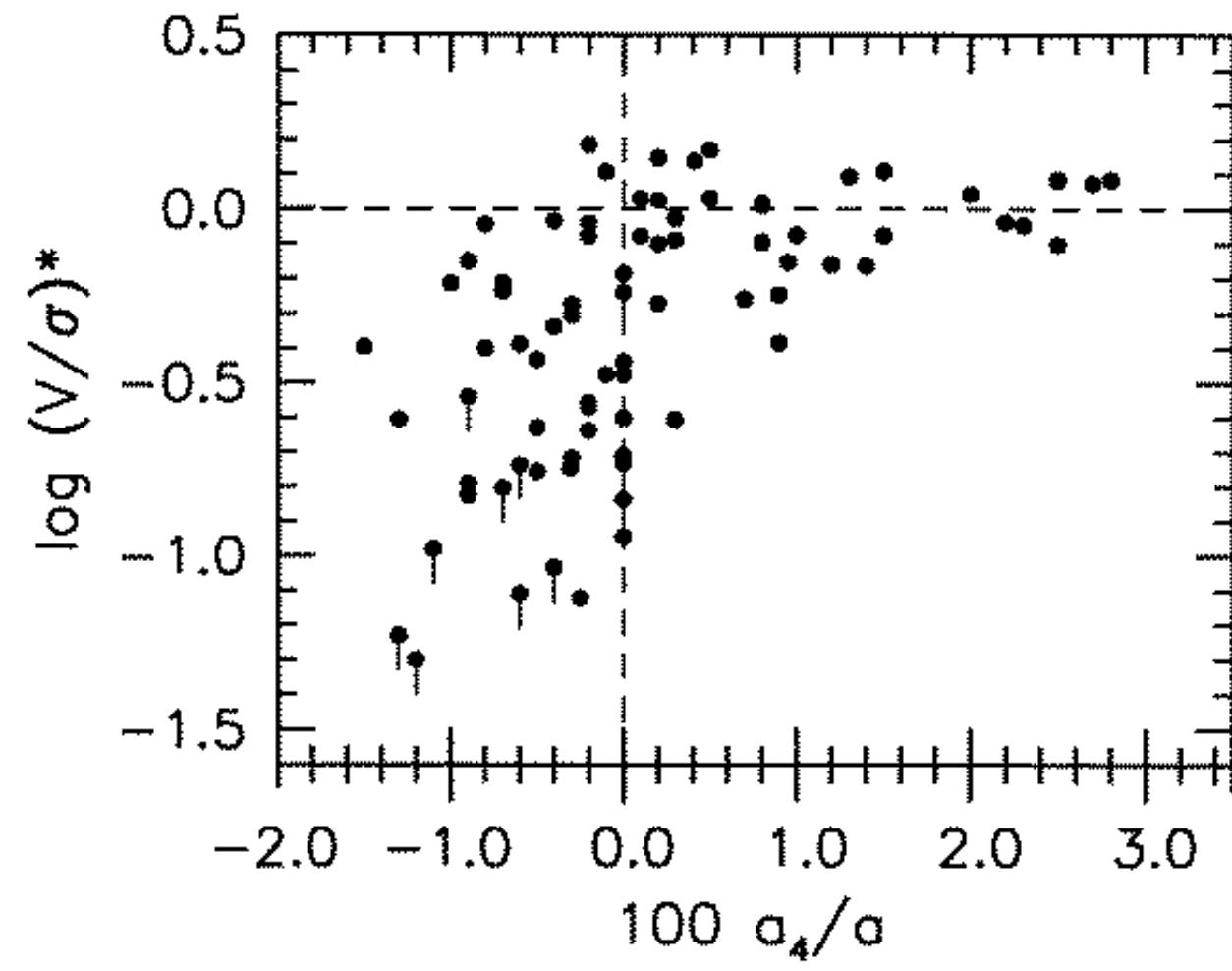
Credit:
Pellegrini Silvia

Disky: NGC4660

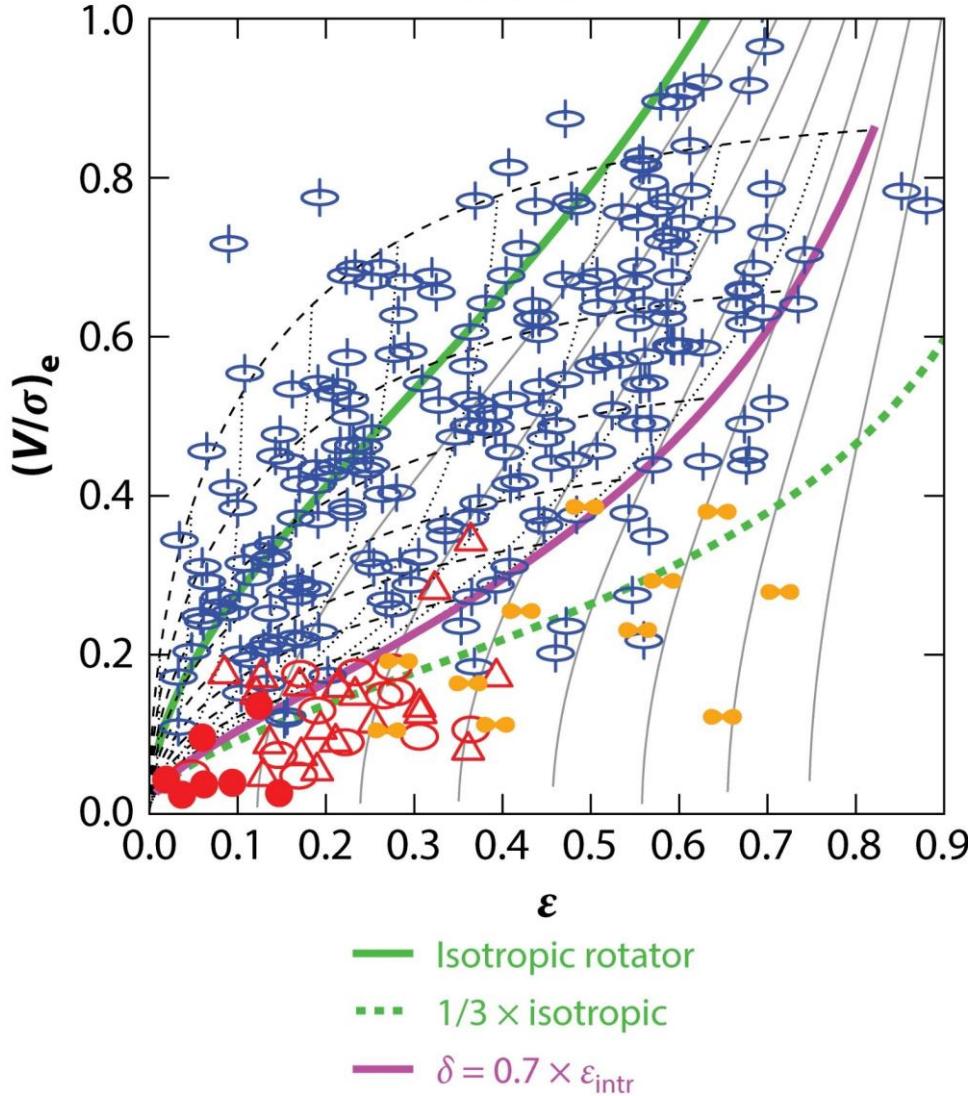


Boxy: NGC5322





ATLAS^{3D}



No disk

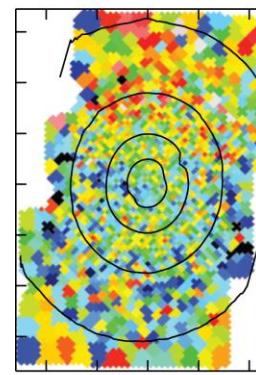
a = ●

b = ○

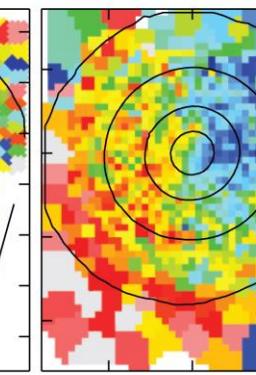
c = ▲

d = ○○

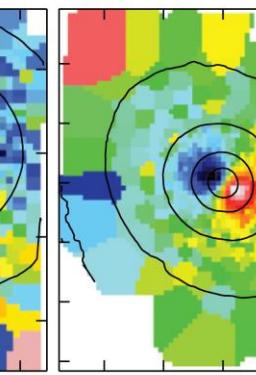
e = ○○○



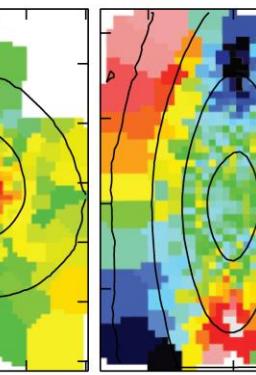
No rotation



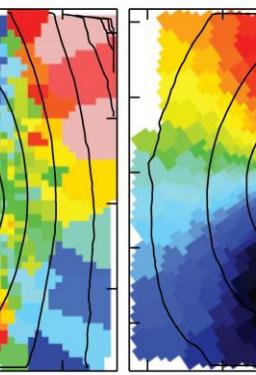
Complex velocity



KDC



Counterrotation



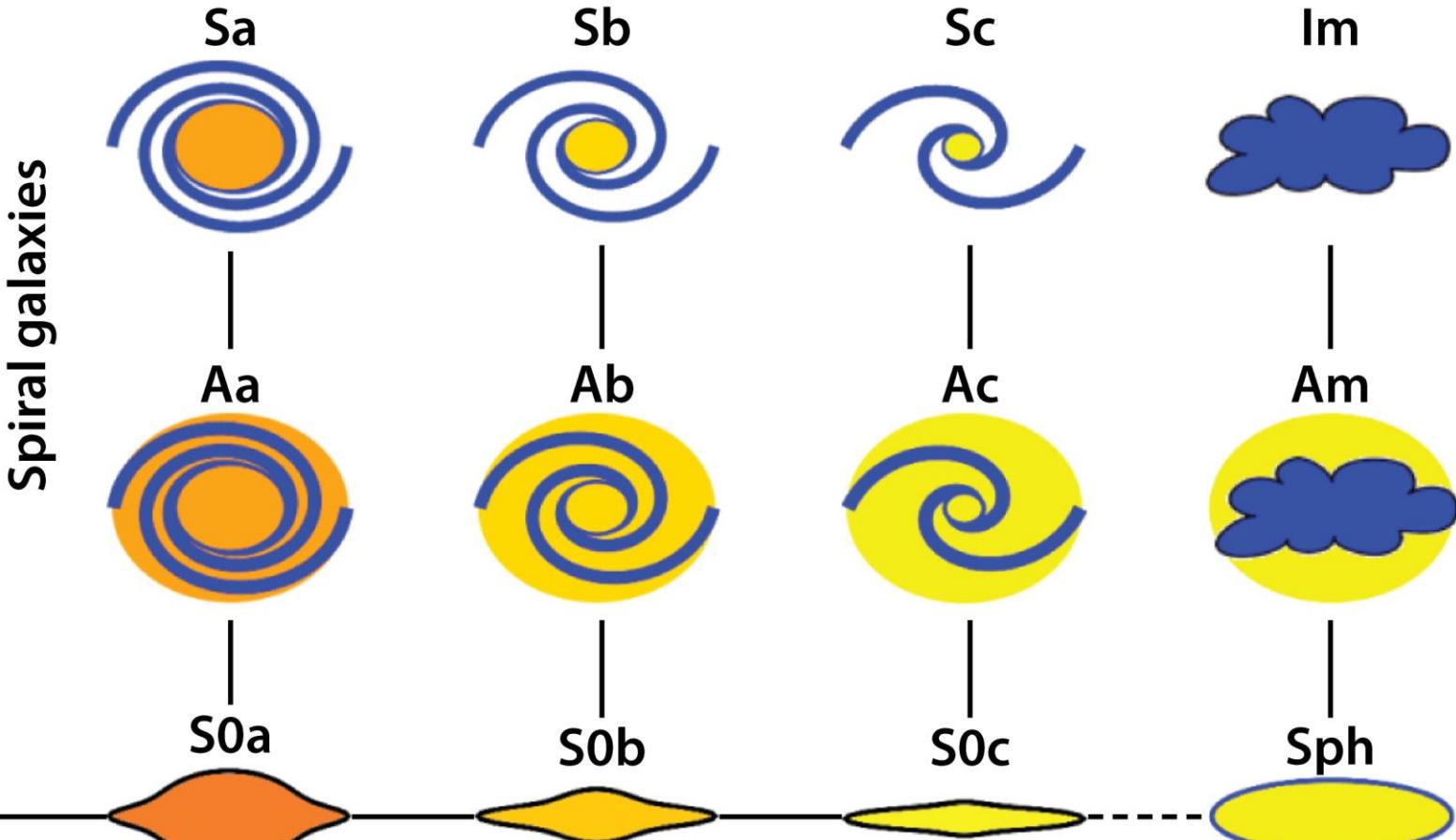
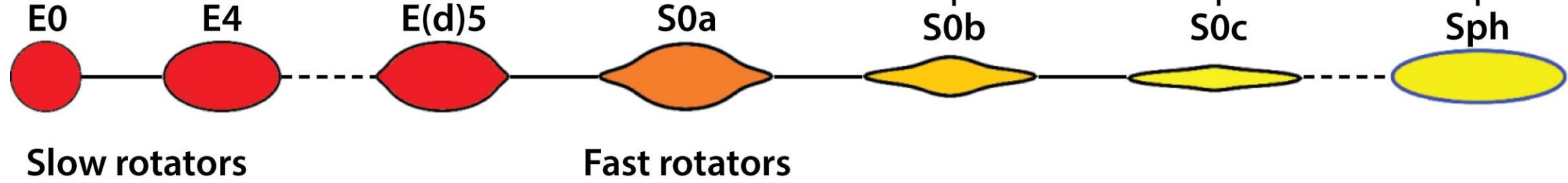
Stellar disk

Nonregular rotators

Regular rotators

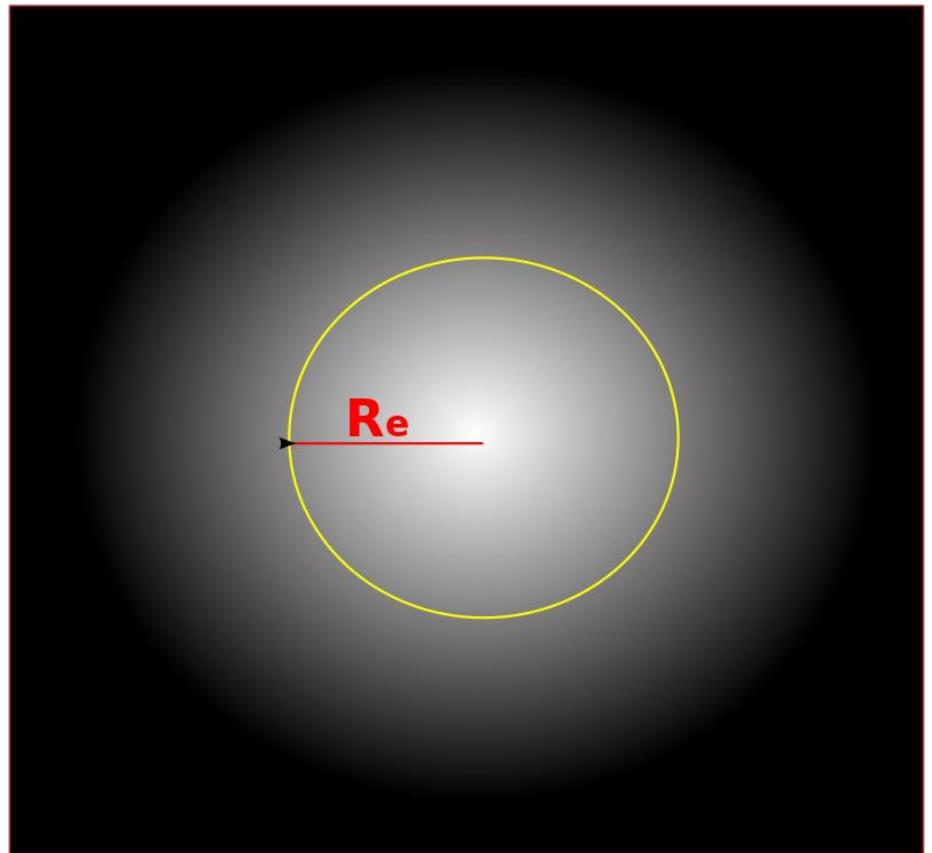
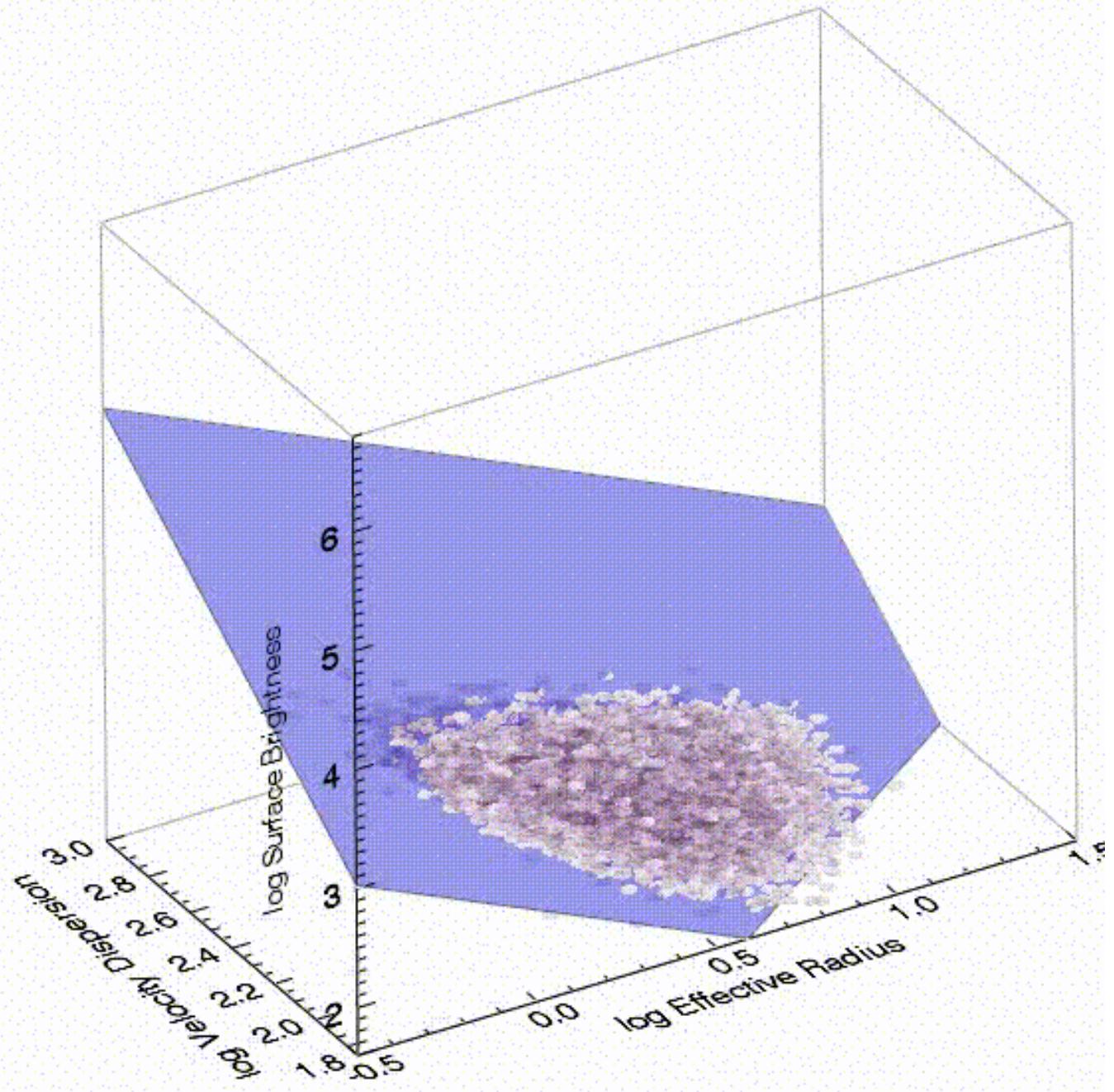
Cappellari M. 2016.

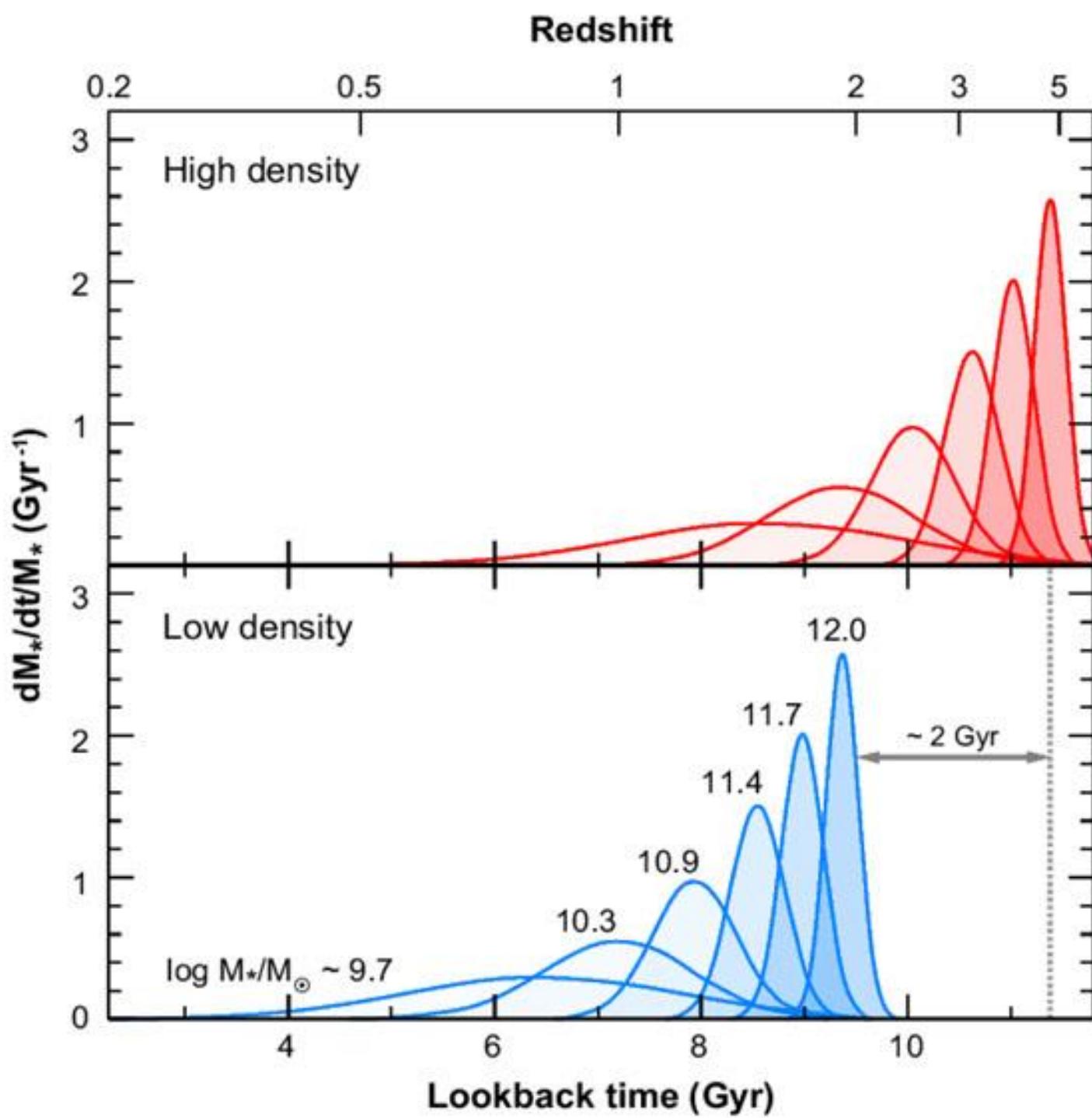
Ann. Rev. Astron. Astrophys. 54:597–665



Cappellari M. 2016.

Annu. Rev. Astron. Astrophys. 54:597–665





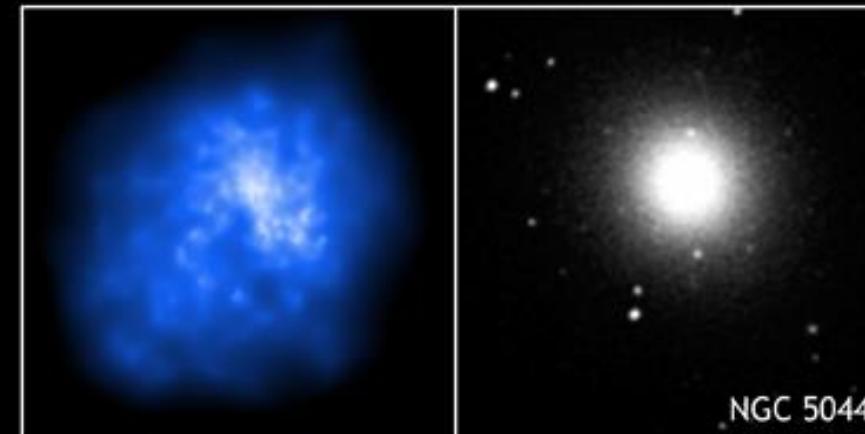
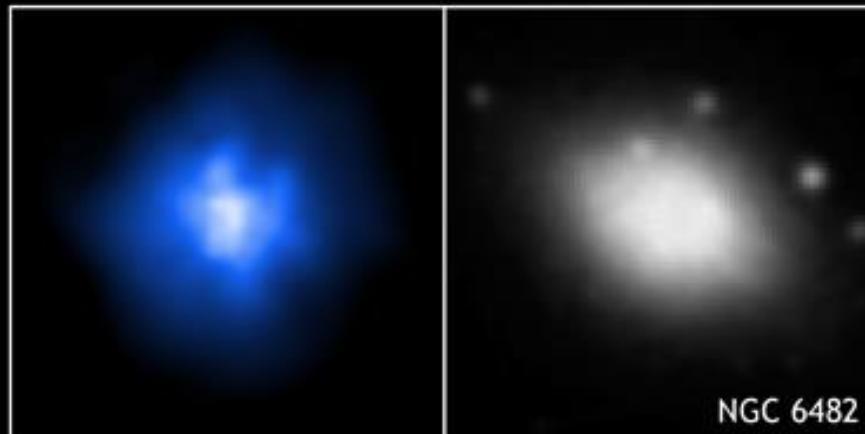
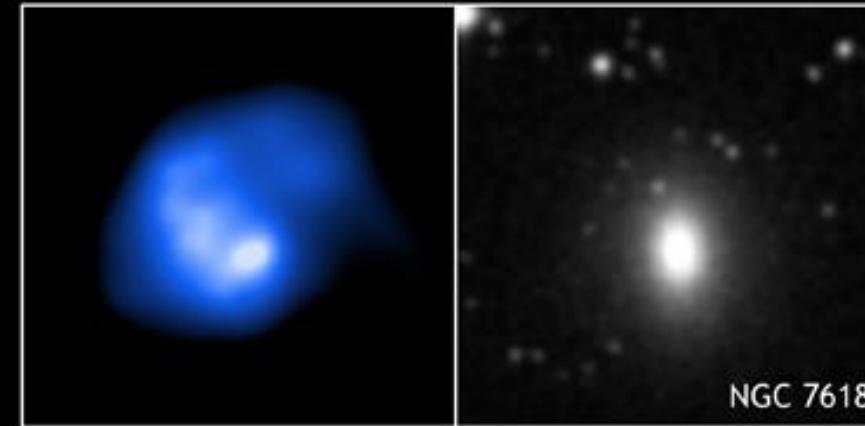
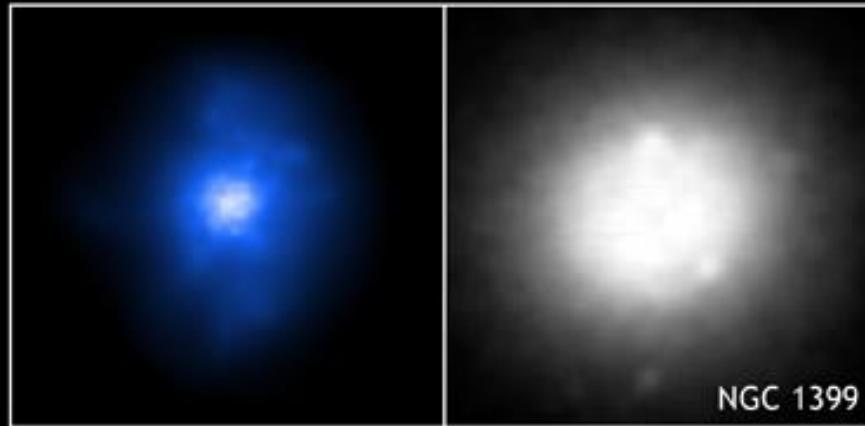
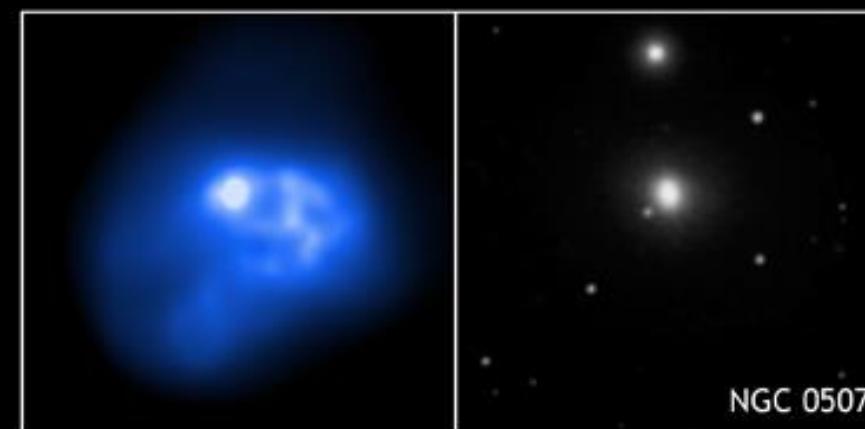
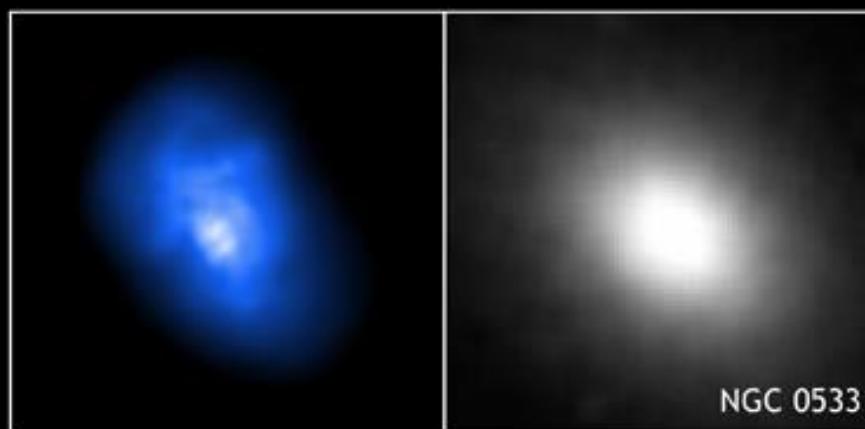
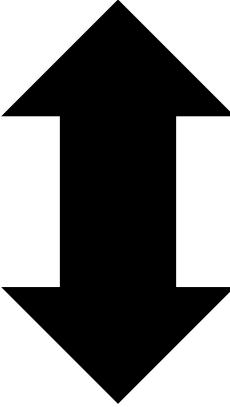




Photo copyright by

$$\frac{dP(r)}{dr} = -\rho(r) \frac{d\phi(r)}{dr} \text{ con } P(r) = \frac{\rho(r)K_B T(r)}{m}$$



$$\frac{d[\rho(r)\sigma^2(r)]}{dr} = -\rho(r) \frac{d\phi(r)}{dr}$$

$$-\frac{d\phi(r)}{dr} = g(r) = -\frac{GM(r)}{r^2} \text{ infatti } F_G = -\frac{GmM(r)}{r^2}$$

$$\left\{ \begin{array}{l} M_T(r) = -\frac{r^2}{G\rho(r)} \frac{d[\rho(r)\sigma^2(r)]}{dr} \\ M_*(r) = 4\pi \int_0^r x^2 \rho(x) dx \end{array} \right.$$

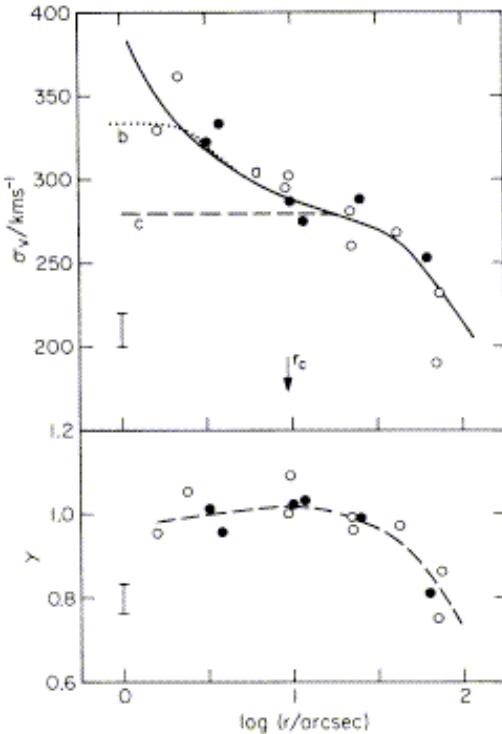


FIG. 4.—Velocity dispersions (σ_v) and the line strengths (γ) for M87. The open circles (\circ) are points W of the nucleus, and the filled circles (\bullet) are E. The core radius, $r_c = 9.6$, of the galaxy is marked. Error bars of length 2σ are given. Curve (a) is the velocity dispersion predicted by the black hole model fitted to the photometric data, (b) is the same model convolved with the seeing disk and slit size for the spectroscopic observations, and (c) is the King model that would prevail if the black hole were absent.

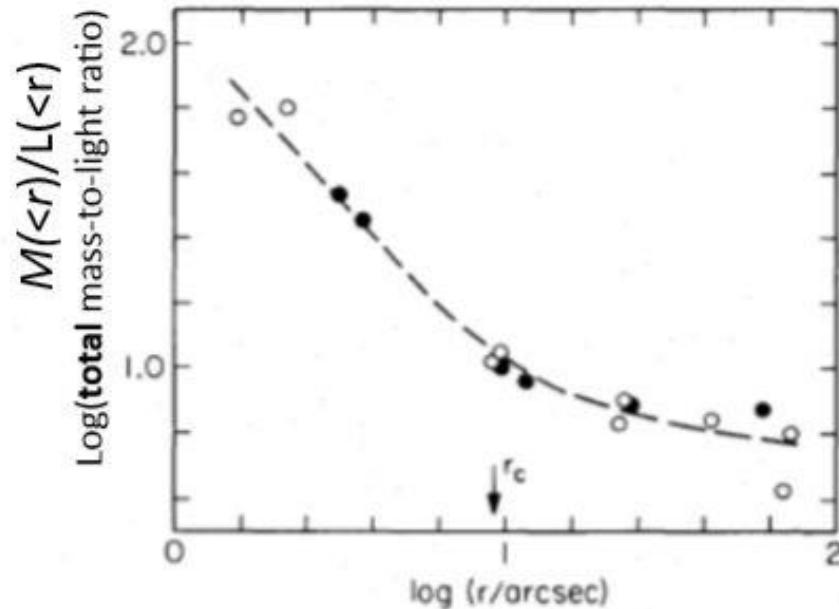


FIG. 10.—Mean mass-luminosity ratio enclosed within radius r [$\langle M/L \rangle(r)$] as a function of r in M87. Open circles (\circ) are to the W of the nucleus, and filled circles (\bullet) are to the E. The core radius of the galaxy, $r_c = 9.6$, is marked.

$M(r)/L(r) \neq \text{const}$

The data require a central $5 \times 10^9 M_\odot$ MBH: the total M/L is consistent with that of the stellar population ($M_*/L \approx \text{constant}$ at the value 6–7) around the core radius r_c , but steeply increases inward

Credit:
Pellegrini Silvia

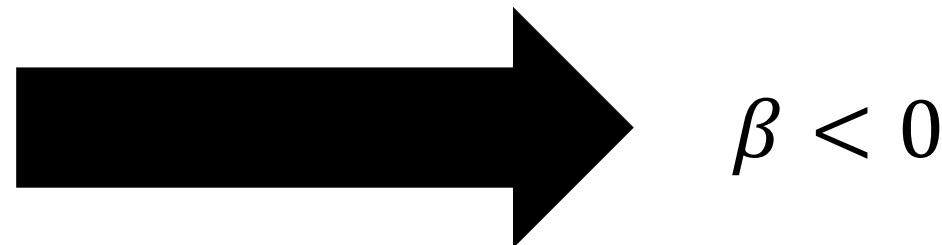
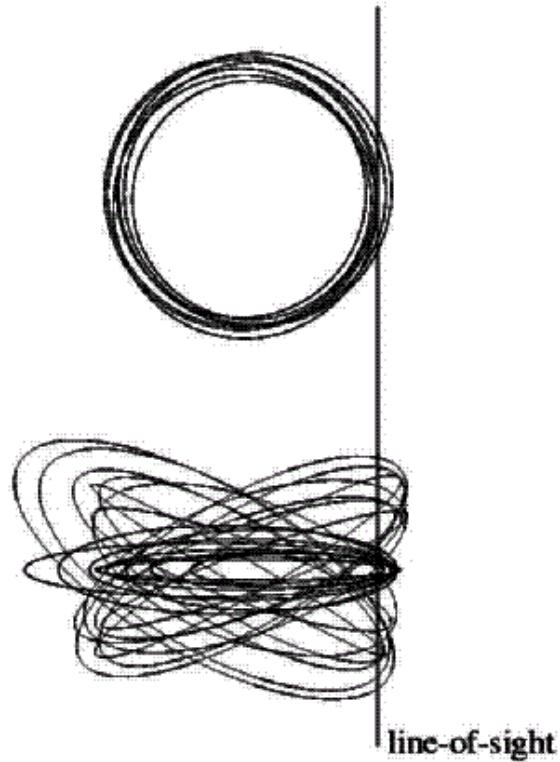
$$\frac{dP(r)}{dr} = -\rho(r) \frac{d\phi(r)}{dr} \text{ con } P(r) = \frac{\rho(r) K_B T(r)}{m}$$



$$\frac{d[\rho(r)\sigma_r^2(r)]}{dr} + \frac{2\beta(r)}{r}\rho(r)\sigma_r^2(r) = -\rho(r)\frac{GM(r)}{r^2}$$

- Assumiamo che tutta la massa sia stellare: $M(r) = M_*(r) \quad \forall r > 0$

Parametro di anisotropia

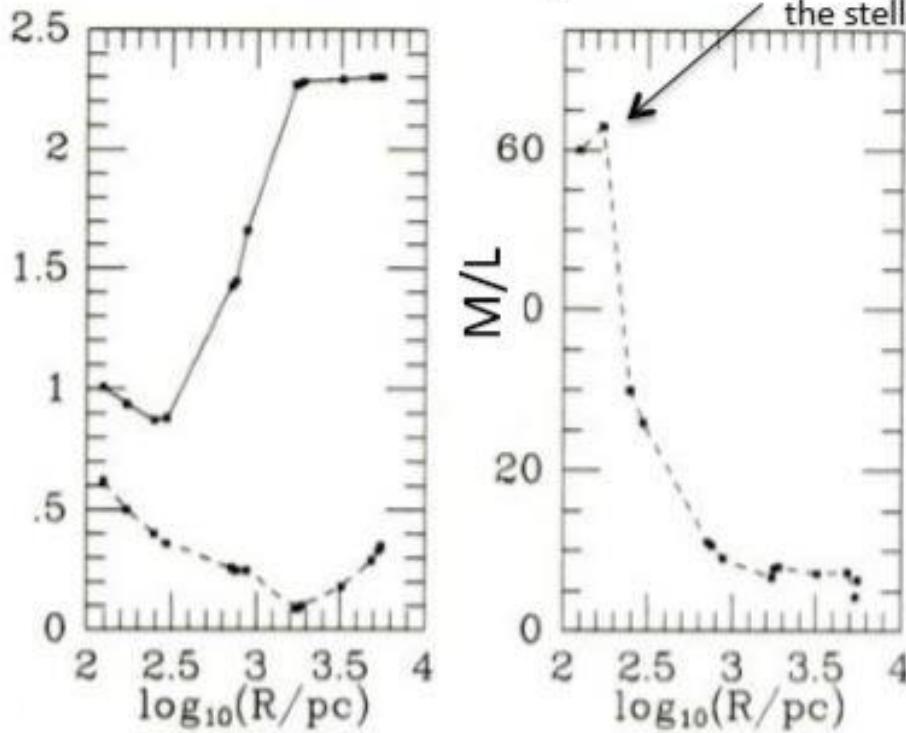


$$\beta < 0$$



$$\beta > 0$$

ISOTROPIC modeling



at the center, M/L becomes
far larger than reasonable for
the stellar population

ANISOTROPIC modeling with a radially varying $\beta(r)$

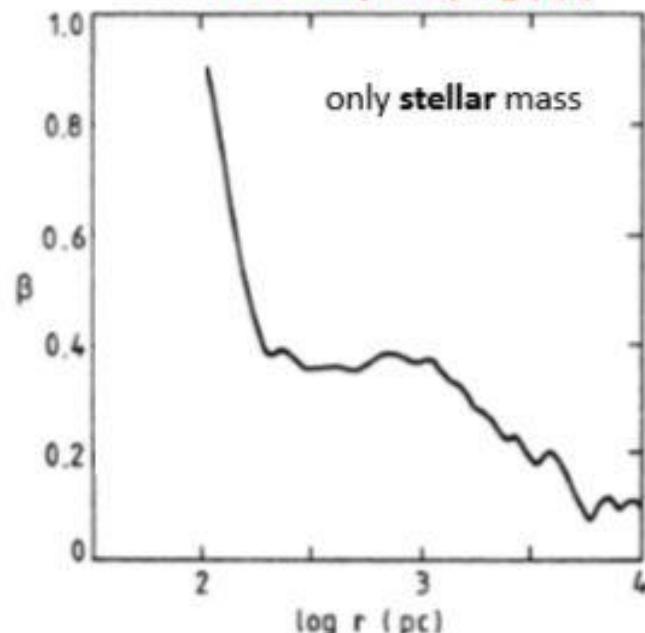


Figure 4-3. Left panel: plots of the logarithmic gradients appearing in equation (4-56) for the giant elliptical galaxy M87: full curve $|d \ln \nu / d \ln r|$; dashed curve $|d \ln \overline{v_r^2} / d \ln r|$. A distance to the galaxy of 16 Mpc has been assumed, and $\overline{v_r^2}$ has been obtained from the observed dispersion under the assumption $\beta = 0$. Right panel: the ratio $M(r)/L(r)$ derived from these data and equation (4-56). (After Sargent et al. 1978.)

Bottom: the radial variation of the anisotropy parameter β in M87 if the observations on which Figure 4-3 is based are to be interpreted in terms of a constant mass-to-light ratio. (After Binney & Mamon 1982.)

$\beta = \text{cost} = 0$ and $M(r)/L(r) \neq \text{cost}$

AND

$\beta(r) > 0$ and $M/L = \text{cost}$

can reproduce equally well the data

BH

NO BH



FINE

Libri consigliati (livello avanzato)

- [Binney and Tremaine,<<Galactic Dynamics>>](#)
- <https://www.annualreviews.org/doi/10.1146/annurev-astro-082214-122432>
- <http://articles.adsabs.harvard.edu/pdf/1978ApJ...221..731S>
- <http://articles.adsabs.harvard.edu//full/1988A%26AS...74..385B/0000391.000.html>
- <http://articles.adsabs.harvard.edu/pdf/1988A%26A...193L...7B>